

Statistical properties for flows with unbounded roof function, including the Lorenz attractor

Péter Bálint ^{*} Ian Melbourne [†]

5 March 2018. Revised 23 June 2018

Abstract

For geometric Lorenz attractors (including the classical Lorenz attractor) we obtain a greatly simplified proof of the central limit theorem which applies also to the more general class of codimension two singular hyperbolic attractors.

We also obtain the functional central limit theorem and moment estimates, as well as iterated versions of these results. A consequence is deterministic homogenisation (convergence to a stochastic differential equation) for fast-slow dynamical systems whenever the fast dynamics is singularly hyperbolic of codimension two.

1 Introduction

The classical Lorenz attractor for the Lorenz equations [24],

$$\dot{x} = 10(y - x), \quad \dot{y} = 28x - y - xz, \quad \dot{z} = xy - \frac{8}{3}z,$$

is a prototypical example of a chaotic dynamical system. There has been much interest in studying rigorously the properties of geometric Lorenz attractors [3, 16] which resemble the numerically-observed classical Lorenz attractor. A computer-assisted proof that the attractor for the classical Lorenz equations is indeed a geometric Lorenz attractor was subsequently provided by [35]. In contrast to Axiom A attractors [34], the Lorenz attractor contains an equilibrium (steady-state for the flow) and cannot be structurally stable, though it is robustly nontrivial and transitive [35]. Moreover, the Lorenz attractor is *singular hyperbolic* [30] and exhibits strong statistical properties

^{*}MTA-BME Stochastics Research Group, Budapest University of Technology and Economics, Egrý József u. 1, H-1111 Budapest, Hungary and Institute of Mathematics, Budapest University of Technology and Economics, H-1111 Egrý József u. 1 Budapest, Hungary. pet@math.bme.hu

[†]Mathematics Institute, University of Warwick, Coventry, CV4 7AL, UK. i.melbourne@warwick.ac.uk

such as expansivity [9], the central limit theorem (CLT) [17], moment estimates [28], and exponential decay of correlations [6].

Recently, [31] (see also previous work of [14]) gave an analytic proof of existence of geometric Lorenz attractors in the extended Lorenz model

$$\dot{x} = y, \quad \dot{y} = -\lambda y + \gamma x(1 - z) - \delta x^3, \quad \dot{z} = -\alpha z + \beta x^2.$$

A consequence of the current paper, in conjunction with [8], is that CLTs and moment estimates hold for these attractors. Exponential decay of correlations remains an open question: the proof in [6] relies on the existence of a smooth stable foliation for the flow, which holds for the classical Lorenz attractor [7] but not for the examples of [14, 31]. (As far as the results on the CLT and moment estimates go, our results for the extended Lorenz model are completely analytic, whereas results for the classical Lorenz attractor rely on the computer-assisted proof in [35].)

More generally, our results apply to all singular hyperbolic attractors for three-dimensional flows. By [30], this incorporates all nontrivial robustly transitive attractors for three-dimensional flows (including nontrivial Axiom A attractors as well as classical and geometric Lorenz attractors). Our results apply also to codimension two singular hyperbolic attractors in arbitrary dimension.

A standard step in the analysis of Lorenz attractors and singular hyperbolic attractors is to consider a suitable Poincaré map with good nonuniform hyperbolicity properties. The return time function (roof function) is generally unbounded due to the presence of steady-states for the flow. Due to this unboundedness, the original proof of the CLT for the classical Lorenz equations [17] is quite complicated, relying on an inducing scheme that involves phase-space exclusion arguments of the type in [12] to remove points that return too quickly to a vicinity of the steady-states.

As we pointed out in an earlier unpublished preprint version of this paper, it is possible to give a much simpler proof of the CLT than the one in [17]. In this paper, we provide the details. As in [17], we prove also the functional version, namely the weak invariance principle (WIP). We also provide moment estimates as advertised in [28, Section 5.3] and our results apply to the general class of singular hyperbolic attractors in [8] (it is not clear that the argument in [17] applies in this generality). In addition, we prove iterated versions of the WIP and moment estimates as advertised in [18, Section 10.3]. By [18, 19] it is therefore possible to prove homogenisation results where fast-slow systems converge to a limiting stochastic differential equation.

Remark 1.1 One version of an earlier preprint of this paper was called “Decay of correlations for flows with unbounded roof function, including the infinite horizon planar periodic Lorentz gas”. The correlation decay aspects of that preprint were incorrect and are subsumed by [10]. The remainder of that preprint is the basis of the current paper.

We note also that a “rearrangement” argument plays a much larger role in the older versions, but is still useful for optimal moment control when the L^p class of the

roof function is restricted, see Section 3 and Remark 4.6(ii). (Rearrangement turns out not to be required for singular hyperbolic attractors where the roof function lies in L^p for all $p < \infty$.)

1.1 Consequences for singular hyperbolic attractors

A specific example to which our main results apply is the classical Lorenz attractor [24, 35] (or more generally, the class of geometric Lorenz attractors [3, 16]). More generally still, we consider singular hyperbolic attractors [30].

Definition 1.2 Let $T_t : \mathbb{R}^n \rightarrow \mathbb{R}^n$ be a C^r flow where $n \geq 3$, $r > 1$. A compact invariant set Λ is *singular hyperbolic* if

- (i) All steady-states in Λ are hyperbolic.
- (ii) The tangent bundle over Λ can be written as a continuous DT_t -invariant sum $T_\Lambda M = E^s \oplus E^{cu}$, and there exist constants $C > 0$, $\lambda \in (0, 1)$ such that for all $x \in \Lambda$, $t \geq 0$,

domination of the splitting: $\|DT_t|E_x^s\| \cdot \|DT_{-t}|E_{T_t x}^{cu}\| \leq C\lambda^t$,

uniform contraction along E^s : $\|DT_t|E_x^s\| \leq C\lambda^t$,

volume expansion along E^{cu} : $|\det(DT_t|E_x^{cu})| \geq C^{-1}\lambda^{-t}$.

Such an invariant set is *codimension two* if $\dim E_x^{cu} = 2$ for $x \in \Lambda$.

A *singular hyperbolic attractor* is a singular hyperbolic set that is also an attractor (attracting and transitive). We recall that Λ is attracting if there exists a positively invariant set $U \subset \mathbb{R}^n$ such that $\Lambda = \bigcap_{t \geq 0} T_t U$, and that Λ is transitive if there exists a dense orbit in Λ .

Remark 1.3 (a) If Λ is a codimension two singular hyperbolic attractor and $x_0 \in \Lambda$ is a steady-state, then x_0 is *Lorenz-like*: the linearized vector field at x_0 restricted to $E_{x_0}^{cu}$ has real eigenvalues λ^s, λ^u satisfying $-\lambda^u < \lambda^s < 0 < \lambda^u$.

(b) By [8, Section 8.1], an attracting codimension two singular hyperbolic attracting set admits a spectral decomposition into a finite disjoint union of singular hyperbolic attractors. Hence, there is no loss of generality in assuming transitivity.

(c) As explained in Section 6, the results for singular hyperbolic attractors described below make use of [8] which relies heavily on the codimension two assumption. Statistical properties for singular hyperbolic attractors with $\dim E_x^{cu} > 2$ remains an interesting open question.

We now describe the main results of this paper in the setting of codimension two singular hyperbolic attractors $\Lambda \subset \mathbb{R}^n$. More general statements for flows with unbounded roof function are given in the body of the paper.

Notation For $a, b \in \mathbb{R}^d$, we define the outer product $a \otimes b = ab^T \in \mathbb{R}^{d \times d}$. For real-valued functions f, g , the integral $\int f dg$ denotes the Itô integral (where defined). Similarly, for vector-valued functions, $\int f \otimes dg$ denotes matrices of Itô integrals. Also $\int f \otimes \circ dg$ denotes matrices of Stratonovich integrals.

Let $\eta \in (0, 1]$. Given $v : M \rightarrow \mathbb{R}^d$, where M is a metric space, define $|v|_\eta = \sup_{x \neq x'} |v(x) - v(x')|/d(x, x')^\eta$ and $\|v\|_\eta = |v|_\infty + |v|_\eta$. Let $C^\eta(M, \mathbb{R}^d)$ be the space of observables $v : M \rightarrow \mathbb{R}^d$ with $\|v\|_\eta < \infty$. When $d = 1$, we write $C^\eta(M)$.

Similarly, for $k \geq 0$ an integer, we write $b \in C^{k, \eta}(\mathbb{R}^d \times M, \mathbb{R}^d)$ if $b = b(x, y)$ is C^k in x uniformly in y , and C^η in y uniformly in x . We write $b \in C^{k+, \eta}(\mathbb{R}^d \times M, \mathbb{R}^d)$ if in addition there exists $\eta' \in (0, 1)$ such that $\sup_y \|\partial b^k / \partial x^k(\cdot, y)\|_{\eta'} < \infty$.

If μ is a specified Borel probability measure on M , then we define $C_0^\eta(M, \mathbb{R}^d) = \{v \in C^\eta(M, \mathbb{R}^d) : \int_M v d\mu = 0\}$.

Let $T_t : \mathbb{R}^n \rightarrow \mathbb{R}^n$ denote a flow with codimension two singular hyperbolic attractor Λ . By [8, 9, 23], there is a unique SRB (Sinai-Ruelle-Bowen) probability measure μ supported on Λ . In particular, μ is a physical measure: there is an open set $U \subset \mathbb{R}^n$ and a subset $U' \subset U$ with $\text{Vol}(U \setminus U') = 0$ such that $\lim_{t \rightarrow \infty} t^{-1} v_t(x_0) = \int_\Lambda v d\mu$ for every continuous function $v : \mathbb{R}^n \rightarrow \mathbb{R}$ and every initial condition $x_0 \in U'$. Here, $v_t = \int_0^t v \circ T_s ds$. (In the case of the classical Lorenz attractor, $U = \mathbb{R}^3$.) Fix $\eta \in (0, 1]$.

Theorem 1.4 (CLT) *Let $v \in C_0^\eta(\Lambda)$. Then the limit $\sigma^2 = \lim_{t \rightarrow \infty} t^{-1} \int_\Lambda v_t^2 d\mu$ exists, and for all $c \in \mathbb{R}$,*

$$\lim_{t \rightarrow \infty} \mu(x_0 \in \Lambda : t^{-1/2} v_t(x_0) \leq c) = (2\pi\sigma^2)^{-1/2} \int_{-\infty}^c e^{-y^2/(2\sigma^2)} dy.$$

(That is, $t^{-1/2} v_t \rightarrow_d N(0, \sigma^2)$ as $t \rightarrow \infty$.) Moreover, there is a closed subspace $S \subset C_0^\eta(\Lambda)$ of infinite codimension such that $\sigma^2 > 0$ for all $v \in C_0^\eta(\Lambda) \setminus S$.

Next, let $W \in C[0, 1]$ denote Brownian motion with variance σ^2 and define $W_n(t) = n^{-1/2} v_{nt}$. We obtain the weak invariance principle:

(WIP) $W_n \rightarrow_w W$ in $C[0, 1]$ as $n \rightarrow \infty$, for $v \in C_0^\eta(\Lambda)$.

Also, we obtain the following moment estimate:

(Moments) For any $p \in (0, \infty)$ there exists $C > 0$ such that $|v_t|_p \leq Ct^{1/2} \|v\|_\eta$ for all $v \in C_0^\eta(\Lambda)$, $t > 0$.

Corollary 1.5 $\lim_{t \rightarrow \infty} \int_\Lambda |t^{-1/2} v_t|^p d\mu = \mathbb{E}|Y|^p$ for all $p \in (0, \infty)$, where Y is normally distributed with mean zero and variance σ^2 .

Proof This is a standard consequence of the CLT and moment estimates, see [28, Lemma 2.1(e)]. ■

Next we consider iterated versions of the WIP and moment estimates. Let $v \in C_0^\eta(\Lambda, \mathbb{R}^d)$ and define $W_n : \Lambda \rightarrow C([0, 1], \mathbb{R}^d)$ as before. Also, define

$$\mathbb{W}_n : \Lambda \rightarrow C([0, 1], \mathbb{R}^{d \times d}), \quad \mathbb{W}_n(t) = \int_0^t W_n \otimes dW_n.$$

(This is a standard Riemann integral since W_n is C^1 in t .)

Theorem 1.6 (Iterated WIP) *The limits*

$$\Sigma = \lim_{n \rightarrow \infty} \int_{\Lambda} W_n(1) \otimes W_n(1) d\mu, \quad E = \lim_{n \rightarrow \infty} \int_{\Lambda} \mathbb{W}_n(1) d\mu,$$

exist and $(W_n, \mathbb{W}_n) \rightarrow_w (W, \mathbb{W})$ in $C([0, 1], \mathbb{R}^d \times \mathbb{R}^{d \times d})$, where W is d -dimensional Brownian motion with covariance Σ and

$$\mathbb{W}(t) = \int_0^t W \otimes dW + Et = \int_0^t W \otimes \circ dW + Dt, \quad D = E - \frac{1}{2}\Sigma.$$

Moreover, there is a closed subspace $S \subset C_0^\eta(\Lambda, \mathbb{R}^d)$ of infinite codimension such that $\det \Sigma > 0$ for all $v \in C_0^\eta(\Lambda, \mathbb{R}^d) \setminus S$.

Remark 1.7 The matrices $\Sigma, D \in \mathbb{R}^{d \times d}$ satisfy $\Sigma^T = \Sigma$ and $D^T = -D$. If in addition the flow T_t is mixing, then

$$\begin{aligned} \Sigma &= \int_0^\infty \int_{\Lambda} (v \otimes (v \circ T_t) + (v \circ T_t) \otimes v) d\mu dt, \\ D &= \frac{1}{2} \int_0^\infty \int_{\Lambda} (v \otimes (v \circ T_t) - (v \circ T_t) \otimes v) d\mu dt. \end{aligned}$$

Next, define $S_t = \int_0^t \int_0^s (v \circ T_r) \otimes (v \circ T_s) dr ds$. We obtain

(Iterated moments) For any $p \in (0, \infty)$ there exists $C > 0$ such that $|S_t|_p \leq Ct \|v\|_\eta^2$ for all $v \in C_0^\eta(\Lambda, \mathbb{R}^d)$, $t > 0$.

Finally, we mention the consequence of these results for deterministic homogenisation. Consider a fast-slow system of the form

$$\dot{x} = a(x, y) + \epsilon^{-1}b(x, y) \quad \dot{y} = \epsilon^{-2}g(y),$$

with $x \in \mathbb{R}^d$ and $y \in \Lambda \subset \mathbb{R}^n$ where Λ is a codimension two singular hyperbolic attractor for the ODE $\dot{y} = g(y)$. We suppose that $x(0) = \xi$ is fixed and that $y(0)$ is chosen randomly from the probability space (Λ, μ) . Apart from the choice of initial condition $y(0)$, the system is completely deterministic.

Theorem 1.8 Suppose that $a \in C^{1+,0}(\mathbb{R}^d \times \Lambda, \mathbb{R}^d)$ and $b \in C_0^{2+,\eta}(\mathbb{R}^d \times \Lambda, \mathbb{R}^d)$ for some $\eta \in (0, 1]$. For $v, w \in C_0^\eta(\Lambda)$, let $\mathfrak{B}(v, w) = \lim_{n \rightarrow \infty} n^{-1} \int_\Lambda \int_0^n \int_0^s v \circ T_r w \circ T_s dr ds d\mu$. Define $\tilde{a} : \mathbb{R}^d \rightarrow \mathbb{R}^d$ and $\sigma : \mathbb{R}^d \rightarrow \mathbb{R}^{d \times d}$,

$$\begin{aligned} \tilde{a}^i(x) &= \int_\Lambda a^i(x, y) d\mu(y) + \sum_{k=1}^d \mathfrak{B}(b^k(x, \cdot), \partial_k b^i(x, \cdot)), \\ (\sigma(x)\sigma(x)^T)^{ij} &= \mathfrak{B}(b^i(x, \cdot), b^j(x, \cdot)) + \mathfrak{B}(b^j(x, \cdot), b^i(x, \cdot)), \end{aligned}$$

for $i, j = 1, \dots, d$.

Then $x_\epsilon \rightarrow_w X$ in $C([0, 1], \mathbb{R}^d)$ where X is the unique solution of the SDE

$$dX = \tilde{a}(X) dt + \sigma(X) dW, \quad X(0) = \xi.$$

Proof We apply [19, Theorem 2.3]. Define the bilinear operator

$$\mathfrak{B} : C_0^\eta(\Lambda) \times C_0^\eta(\Lambda) \rightarrow \mathbb{R}, \quad \mathfrak{B}(u, u') = \lim_{n \rightarrow \infty} \int_\Lambda \int_0^1 W_{u,n} \otimes dW_{u',n} d\mu,$$

where $W_{u,n} = n^{-1/2}u_{nt}$. By Theorem 1.6, for every $m \geq 1$ and $v \in C_0^\eta(\Lambda, \mathbb{R}^m)$, we have that $(W_n, \mathbb{W}_n) \rightarrow_w (W, \mathbb{W})$ in $C([0, 1], \mathbb{R}^m \times \mathbb{R}^{m \times m})$, and hence in the sense of finite-dimensional distributions, where W is a Brownian motion and $\mathbb{W}(t) = \int_0^t W \otimes dW + Et$. Moreover,

$$E^{ij} = \lim_{n \rightarrow \infty} \frac{1}{n} \int_\Lambda \mathbb{W}_n(1)^{ij} d\mu = \lim_{n \rightarrow \infty} \frac{1}{n} \int_\Lambda \int_0^1 W_n^i dW_n^j d\mu = \mathfrak{B}(v^i, v^j),$$

for $i, j = 1, \dots, m$. This is Assumption 2.1 in [19]. Also, Assumption 2.2 in [19] follows from the moment estimates $|v_t|_p \ll t^{1/2}$ and $|S_t|_p \ll t$ which hold for all p and all $v \in C_0^\eta(\Lambda, \mathbb{R}^d)$. Hence, the result follows from [19, Theorem 2.3]. \blacksquare

Remark 1.9 The special case $a(x, y) = a(x)$ and $b(x, y) = b(x)v(y)$ with $a \in C^{1+}(\mathbb{R}^d, \mathbb{R}^d)$, $b \in C^{2+}(\mathbb{R}^d, \mathbb{R}^{d \times e})$, $v \in C_0^\eta(\Lambda, \mathbb{R}^e)$, is considered in [18]. For such fast-slow systems, the limiting SDE takes the form

$$dX = \left\{ a(X) + \sum_{\alpha, \beta, \gamma=1}^d D^{\beta\gamma} \partial_\alpha b^\beta(X) b^{\alpha\gamma}(X) \right\} dt + b(X) \circ dW, \quad X(0) = \xi,$$

where D is as in Theorem 1.6. (Here $b^\beta(X)$ is the β 'th column of $b(X)$.)

The remainder of this paper is organised as follows. In Section 2, we prove the iterated WIP and moment estimates for nonuniformly expanding systems with unbounded roof function. In Section 3, we use a rearrangement idea to obtain sharper moment estimates. Sections 4 and 5 contain the iterated WIP and moment estimates

for two classes of nonuniformly hyperbolic flows; one where there is exponential contraction along stable leaves for the flow, and one where there is a horseshoe structure as in [36]. Finally, in Section 6 we verify the hypotheses in Section 4 for singular hyperbolic attractors. As a consequence of this, we obtain all of the results described in the introduction.

We end this introduction with some basic material on suspensions and inducing.

1.2 Preliminaries on suspensions and inducing

Suspensions Let $T_t : M \rightarrow M$ be a semiflow defined on a measurable space M . (We assume that $(x, t) \mapsto T_t x$ is measurable from $M \times [0, \infty) \rightarrow M$.) Let $X \subset M$ be a measurable subset with probability measure μ_X . Suppose that $f : X \rightarrow X$ is a measure-preserving transformation and $h : X \rightarrow \mathbb{R}^+$ is an integrable function such that $T_{h(x)}x = fx$ for all $x \in X$.

Define the suspension $X^h = \{(x, u) \in X \times [0, \infty) : u \in [0, h(x)]\} / \sim$ where $(x, h(x)) \sim (fx, 0)$. The suspension semiflow $f_t : X^h \rightarrow X^h$ is given by $f_t(x, u) = (x, u + t)$ computed modulo identifications. We obtain an f_t -invariant probability measure on X^h given by $\mu^h = (\mu_X \times \text{Lebesgue}) / \bar{h}$ where $\bar{h} = \int_X h d\mu_X$.

Define $\pi : X^h \rightarrow M$ by $\pi(x, u) = T_u x$. This defines a measurable semiconjugacy from f_t to T_t and $\mu_M = \pi_* \mu^h$ is a T_t -invariant probability measure on M .

Inducing Let $f : X \rightarrow X$ be a measurable map defined on a measurable space X . Let $Y \subset X$ be a measurable subset with probability measure μ_Y . Suppose that $F : Y \rightarrow Y$ is a measure-preserving transformation and $\tau : Y \rightarrow \mathbb{Z}^+$ is an integrable function such that $f^{\tau(y)}y = Fy$ for all $y \in Y$.

Define the tower $Y^\tau = \{(y, \ell) \in Y \times \mathbb{Z} : 0 \leq \ell \leq \tau(y) - 1\}$. The tower map $\hat{f} : Y^\tau \rightarrow Y^\tau$ is given by $\hat{f}(y, \ell) = \begin{cases} (y, \ell + 1) & \ell \leq \tau(y) - 2 \\ (Fy, 0) & \ell = \tau(y) - 1 \end{cases}$. We obtain an \hat{f} -invariant probability measure on Y^τ given by $\mu^\tau = (\mu_Y \times \text{counting}) / \bar{\tau}$ where $\bar{\tau} = \int_Y \tau d\mu_Y$.

Define $\pi : Y^\tau \rightarrow X$ by $\pi(y, \ell) = f^\ell y$. This defines a measurable semiconjugacy from \hat{f} to f and $\mu_X = \pi_* \mu^\tau$ is an f -invariant probability measure on X .

Induced roof functions Often these situations are combined, so we have $T_t : M \rightarrow M$, $f : X \rightarrow X$ and $F : Y \rightarrow Y$ together with $h : X \rightarrow \mathbb{R}^+$ and $\tau : Y \rightarrow \mathbb{Z}^+$ such that $T_h = f$ and $f^\tau = F$.

If μ_Y is an F -invariant probability measure on Y and τ is integrable, then we can define μ_X on X as above. If in addition h is integrable, then we can define μ_M on M as above. It is convenient to define the *induced roof function*

$$H : Y \rightarrow \mathbb{R}^+, \quad H(y) = \sum_{\ell=0}^{\tau(y)-1} h(f^\ell y). \quad (1.1)$$

Then H is integrable (indeed $\int_Y H d\mu_Y = \bar{h}\bar{\tau}$) and we define the suspension semiflow $F_t : Y^H \rightarrow Y^H$ with invariant probability measure $\mu^H = (\mu_Y \times \text{Lebesgue}) / \bar{H}$ where

$\bar{H} = \int_Y H d\mu_Y$. Define the semiconjugacy $\pi : Y^H \rightarrow M$ given by $\pi(y, u) = T_u y$. It is easily checked that $\pi_* \mu^H$ is again μ_M .

More notation We use “big O” and \ll notation interchangeably. For example, $a_n = O(b_n)$ or $a_n \ll b_n$ means that there is a constant $C > 0$ such that $a_n \leq Cb_n$ for all $n \geq 1$.

2 Nonuniformly expanding semiflows

Let (M, d) be a bounded metric space with Borel subsets $Y \subset X \subset M$. Suppose that $T_t : M \rightarrow M$, $h : X \rightarrow \mathbb{R}^+$, $f = T_h : X \rightarrow X$ are as in Subsection 1.2. Also, let $\tau : Y \rightarrow \mathbb{Z}^+$ and $F = f^\tau : Y \rightarrow Y$ be as in Subsection 1.2.

Now let μ_0 be a finite Borel measure on X with $\mu_0(Y) > 0$. We suppose that f is a nonsingular transformation for which μ_0 is ergodic, and that $\tau \in L^1(Y)$. Moreover, we suppose that there is an at most countable measurable partition $\{Y_j\}$ such that τ is constant on partition elements, and there exist constants $\lambda > 1$, $C \geq 1$, $\eta \in (0, 1]$ such that for all $j \geq 1$,

- (1) F restricts to a (measure-theoretic) bijection from Y_j onto Y .
- (2) $d(Fy, Fy') \geq \lambda d(y, y')$ for all $y, y' \in Y_j$.
- (3) $\zeta_0 = \frac{d\mu_0|_Y}{d\mu_0|_{Y \circ F}}$ satisfies $|\log \zeta_0(y) - \log \zeta_0(y')| \leq Cd(Fy, Fy')^\eta$ for all $y, y' \in Y_j$.

There is a unique ergodic F -invariant absolutely continuous probability measure μ_Y on Y . Moreover, $d\mu_Y/d\mu_0 \in L^\infty(Y)$ and $F : Y \rightarrow Y$ is a (full branch) Gibbs-Markov map. Standard references for material on Gibbs-Markov maps are [1, Chapter 4] and [2].

Starting from μ_Y and using that τ is integrable, we can construct an ergodic f -invariant probability measure μ_X on X as in Section 1.2. We assume that $\int_X h d\mu_X$ is integrable, leading to an ergodic T_t -invariant probability measure μ_M on M .

For $\theta \in (0, 1)$, define the symbolic metric $d_\theta(y, y') = \theta^{s(y, y')}$ where the *separation time* $s(y, y')$ is the least integer $n \geq 0$ such that $F^n y$ and $F^n y'$ lie in distinct partition elements Y_j . A function $v : Y \rightarrow \mathbb{R}$ is d_θ -Lipschitz if $|v|_\theta = \sup_{y \neq y'} |v(y) - v(y')|/d_\theta(y, y')$ is finite. Let $F_\theta(Y)$ be the Banach space of Lipschitz functions with norm $\|v\|_\theta = |v|_\infty + |v|_\theta$. More generally (and with a slight abuse of notation), we say that a function $v : Y \rightarrow \mathbb{R}$ is *piecewise Lipschitz* if $|1_{Y_j} v|_\theta = \sup_{y, y' \in Y_j, y \neq y'} |v(y) - v(y')|/d_\theta(y, y')$ is finite for all j .

Let $P : L^1(Y) \rightarrow L^1(Y)$ denote the transfer operator corresponding to $F : Y \rightarrow Y$ (so $\int_Y P v w d\mu_Y = \int_Y v w \circ F d\mu_Y$ for all $v \in L^1(Y)$, $w \in L^\infty(Y)$).

Proposition 2.1 *There exists $\theta, \gamma \in (0, 1)$, $C > 0$, such that*

- (a) $\|P^n v\|_\theta \leq C\gamma^n \|v\|_\theta$ for all $v \in F_\theta(Y)$ with $\int_Y v d\mu_Y = 0$ and all $n \geq 1$.
- (b) If $g \in L^1(Y)$ such that $|1_{Y_j} v|_\infty \leq \inf_{Y_j} g$ and $|1_{Y_j} v|_\theta \leq \inf_{Y_j} g$ for all $j \geq 1$, then $\|Pv\|_\theta \leq C|g|_1$.

Proof Part (a) is proved for instance in [1, 2]. Part (b) is a standard argument (eg. [26, Lemma 2.2]). \blacksquare

Define the induced roof function $H : Y \rightarrow \mathbb{R}^+$ as in (1.1). We assume that H is piecewise d_θ -Lipschitz for some $\theta \in (0, 1)$. Moreover, we suppose that there is a constant $C \geq 1$ such that

$$|1_{Y_j} H|_\theta \leq C \inf_{Y_j} H \quad \text{for } j \geq 1, \quad (2.1)$$

and

$$d(T_t y, T_t y') \leq C\theta^{s(y, y')} \text{ for } y, y' \in Y_j, j \geq 1, t \in [0, H(y)] \cap [0, H(y')]. \quad (2.2)$$

2.1 Martingale-coboundary decomposition

Let $v : M \rightarrow \mathbb{R}^d$ and define the induced observable

$$V : Y \rightarrow \mathbb{R}^d, \quad V(y) = \int_0^{H(y)} v(T_u y) du.$$

Proposition 2.2 *Let $\theta_1 = \theta^\eta$. There is a constant $C > 0$ such that*

$$|1_{Y_j} V|_\infty \leq C|v|_\infty \inf_{Y_j} H, \quad |1_{Y_j} V|_{\theta_1} \leq C\|v\|_\eta \inf_{Y_j} H \quad \text{for } j \geq 1,$$

for all $v \in C^\eta(M, \mathbb{R}^d)$.

Proof By (2.1), $|1_{Y_j} V|_\infty \leq |v|_\infty |1_{Y_j} H|_\infty \ll |v|_\infty \inf_{Y_j} H$. Let $y, y' \in Y_j$ and suppose without loss that $H(y) \geq H(y')$. By (2.2), $|v(T_u y) - v(T_u y')| \leq |v|_\eta d(T_u y, T_u y')^\eta \leq C|v|_\eta \theta^{\eta s(y, y')} = C|v|_\eta d_{\theta_1}(y, y')$ for $u \in [0, H(y')]$. Hence by (2.1),

$$\begin{aligned} |V(y) - V(y')| &\leq \int_{H(y')}^{H(y)} |v(T_u y)| du + \int_0^{H(y')} |v(T_u y) - v(T_u y')| du \\ &\leq |v|_\infty (H(y) - H(y')) + C|v|_\eta H(y') d_{\theta_1}(y, y') \ll \|v\|_\eta \inf_{Y_j} H d_{\theta_1}(y, y'), \end{aligned}$$

as required. \blacksquare

Let U be the Koopman operator $Uv = v \circ F$. It is standard that $PU = I$ and $UP = \mathbb{E}(\cdot | F^{-1}\mathcal{B})$ where \mathcal{B} denotes the underlying σ -algebra on Y .

Proposition 2.3 *Suppose that $H \in L^p(Y)$ where $p \geq 1$. Let $v \in C_0^\eta(M, \mathbb{R}^d)$. Then there exists $m \in L^p(Y, \mathbb{R}^d)$, $\chi \in L^\infty(Y, \mathbb{R}^d)$ such that $V = m + \chi \circ F - \chi$, $m \in \ker P$.*

Moreover, there exists a constant $C > 0$ such that $|m|_p \leq C\|v\|_\eta$ and $|\chi|_\infty \leq C\|v\|_\eta$ for all $v \in C_0^\eta(M, \mathbb{R}^d)$.

Proof By Proposition 2.2, $V \in L^p(Y, \mathbb{R}^d)$ with $|V|_p \leq |v|_\infty |H|_p$.

By Proposition 2.1(b) and Proposition 2.2, $\|PV\|_{\theta_1} \ll \|v\|_\eta |H|_1 \ll \|v\|_\eta$. By Proposition 2.1(a), $\|P^n V\|_{\theta_1} \ll \gamma^n \|PV\|_{\theta_1}$ for $n \geq 1$. In particular, $\chi = \sum_{n=1}^{\infty} P^n V$ converges in $L^\infty(Y)$, and we obtain $V = m + \chi \circ F - \chi$ with $|\chi|_\infty \ll \|v\|_\eta$ and $|m|_p \leq |V|_p + 2|\chi|_\infty \ll \|v\|_\eta$. Finally, $Pm = PV - \chi + P\chi = 0$. \blacksquare

Let $V_n = \sum_{j=0}^{n-1} V \circ F^j$.

Corollary 2.4 *Suppose that $H \in L^p(Y)$ where $p \geq 2$. There is a constant $C > 0$ such that $|\max_{1 \leq k \leq n} |V_k|_p| \leq Cn^{1/2} \|v\|_\eta$ for all $v \in C_0^\eta(M, \mathbb{R}^d)$, $n \geq 1$.*

Proof By Proposition 2.3, $\mathbb{E}(m|F^{-1}\mathcal{B}) = UPm = 0$. Hence $\{m \circ F^n, n \geq 0\}$ is a sequence of reversal martingale differences with respect to the sequence of σ -algebras $\{F^{-n}\mathcal{B}, n \geq 0\}$. Let $m_n = \sum_{j=0}^{n-1} m \circ F^j$. Since $p \geq 2$, it follows from Burkholder's inequality [13, Theorem 3.2 combined with (1.4)] that

$$|\max_{1 \leq k \leq n} |m_k|_p| \leq \frac{p}{p-1} |m_n|_p \ll \left| \left(\sum_{j=1}^n m^2 \circ F^j \right)^{1/2} \right|_p \leq n^{1/2} |m|_p \ll n^{1/2} \|v\|_\eta.$$

Since $V_n = m_n + \chi \circ F^n - \chi$ and $\chi \in L^\infty$, we obtain $|\max_{1 \leq k \leq n} |V_k|_p| \leq |\max_{1 \leq k \leq n} |m_k|_p| + 2|\chi|_\infty \ll n^{1/2} \|v\|_\eta$. \blacksquare

2.2 Central limit theorems

Let $v \in C^\eta(M, \mathbb{R}^d)$ and define

$$W_n(t) = n^{-1/2} \int_0^{nt} v \circ T_s ds, \quad \mathbb{W}_n(t) = n^{-1} \int_0^{nt} \int_0^s (v \circ T_r) \otimes (v \circ T_s) dr ds.$$

Also, define $\hat{v}, \psi : Y^H \rightarrow \mathbb{R}^d$,

$$\hat{v}(y, u) = v(T_u y), \quad \psi(y, u) = \int_0^u v(T_s y) ds.$$

Theorem 2.5 *Suppose that $H \in L^2(Y)$. Let $v \in C_0^\eta(M, \mathbb{R}^d)$. Then there exist $\Sigma, D \in \mathbb{R}^{d \times d}$ with $\Sigma^T = \Sigma$ and $D^T = -D$ such that $(W_n, \mathbb{W}_n) \rightarrow_w (W, \mathbb{W})$ in $C([0, \infty), \mathbb{R}^d \times \mathbb{R}^{d \times d})$ on (M, μ_M) , where W is a d -dimensional Brownian motion with covariance Σ and $\mathbb{W}(t) = \int_0^t W \otimes \circ dW + Dt$.*

Proof There are two main steps where we work with (i) the induced observable $V : Y \rightarrow \mathbb{R}^d$, (ii) the observable $v : M \rightarrow \mathbb{R}^d$. (This is more direct than the procedure indicated in [18, Section 10.3].)

Define

$$W_n^Y(t) = n^{-1/2} \sum_{0 \leq j \leq [nt]-1} V \circ F^j, \quad \mathbb{W}_n^Y(t) = n^{-1} \sum_{0 \leq i < j \leq [nt]-1} (V \circ F^i) \otimes (V \circ F^j).$$

Recall that the induced map $F : Y \rightarrow Y$ is mixing. Using the decomposition in Proposition 2.3 with $p = 2$, it follows from [18, Theorem 4.3] that $(W_n^Y, \mathbb{W}_n^Y) \rightarrow (W^Y, \mathbb{W}^Y)$ in $D([0, \infty), \mathbb{R}^d \times \mathbb{R}^{d \times d})$ on (Y, μ_Y) where W^Y is a d -dimensional Brownian motion with some covariance matrix $\Sigma_Y \in \mathbb{R}^{d \times d}$ and

$$\mathbb{W}^Y(t) = \int_0^t W^Y \otimes dW^Y + E_Y t \quad \text{where} \quad E_Y \in \mathbb{R}^{d \times d}.$$

This completes step (i).

Next, we claim that

- (a) $V \in L^2(Y, \mathbb{R}^d)$, $|\psi| |\hat{v}| \in L^1(Y^H)$.
- (b) $n^{-1/2} \sup_{t \in [0, T]} |\psi \circ f_{nt}| \rightarrow_w 0$.
- (c) $\lim_{n \rightarrow \infty} n^{-1} \left| \max_{1 \leq k \leq n} \left| \sum_{1 \leq j \leq k} V \circ F^j \right| \right|_2 = 0$.

To verify the claim, we proceed as in [18, Remark 6.2]. First, $H \in L^2(Y)$ so $V \in L^2(Y, \mathbb{R}^d)$ by Proposition 2.2. Also $|\psi(y, u)| \leq |v|_\infty H(y)$, so $\int_{Y^H} |\psi| d\mu^H \leq |v|_\infty \bar{H}^{-1} \int_Y H^2 d\mu < \infty$. Hence (a) is satisfied. Item (c) is immediate by Corollary 2.4. To prove (b), we argue as in the proof of [18, Proposition 6.6(b)]. Define $g : Y \rightarrow \mathbb{R}$, $g(y) = \int_0^{H(y)} |v(T_s y)| ds$. Then $g(y) \leq |v|_\infty H(y)$ so $g^2 \in L^1(Y)$. By the ergodic theorem, $n^{-1/2} g \circ F^n \rightarrow 0$ a.e. Define the lap number $N(t) : Y^H \rightarrow \mathbb{N}$ to be the largest integer $n \geq 0$ such that $u + \sum_{j=0}^{n-1} H(F^j y) \leq t$. Then $N(nt)/n \rightarrow \bar{H}^{-1}$ a.e. and so $n^{-1/2} g \circ F^{N(nt)} = (N(nt)/n)^{1/2} N(nt)^{-1/2} g \circ F^{N(nt)} \rightarrow 0$ a.e. Also, $|\psi(y, u)| \leq g(y)$, so

$$|\psi \circ f_t(y, u)| = \left| \psi(F^{N(t)(y, u)} y, u + t - \sum_{j=0}^{N(t)(y, u)-1} H(F^j y)) \right| \leq g(F^{N(t)(y, u)} y).$$

Hence $n^{-1/2} \psi \circ f_{nt} \rightarrow 0$ a.e. and (b) follows. This completes the verification of the claim.

Since the claim holds, the hypotheses of [18, Theorem 6.1] are satisfied (with V, ψ, \hat{v} here playing the roles of \tilde{v}, H, v in [18, Theorem 6.1]). It follows from [18, Theorem 6.1] that the iterated WIP for V on Y implies an iterated WIP for \hat{v} on Y^H and hence, via the semiconjugacy $\pi : Y^H \rightarrow M$, an iterated WIP for v on M . We conclude that $(W_n, \mathbb{W}_n) \rightarrow_w (W, \mathbb{W})$ in $C([0, \infty), \mathbb{R}^d \times \mathbb{R}^{d \times d})$ on (M, μ_M) , where

$$W = \bar{H}^{-1/2} W^Y, \quad \mathbb{W}(t) = \int_0^t W \otimes dW + Et, \quad E = \bar{H}^{-1} E_Y + \int_{Y^H} \psi \otimes \hat{v} d\mu^H.$$

The Stratonovich correction gives $D = E - \frac{1}{2} \Sigma$ skew-symmetric as in the proof of [18, Corollary 8.1]. ■

Remark 2.6 By [18, Theorem 6.1],

$$\begin{aligned}\Sigma &= \bar{H}^{-1} \left(\int_Y V \otimes V d\mu_Y + \sum_{n=1}^{\infty} \int_Y \{V \otimes (V \circ F^n) + (V \circ F^n) \otimes V\} d\mu_Y \right), \\ 2D &= \bar{H}^{-1} \sum_{n=1}^{\infty} \int_Y \{V \otimes (V \circ F^n) - (V \circ F^n) \otimes V\} d\mu_Y + \int_{Y^H} (\psi \otimes \hat{v} - \hat{v} \otimes \psi) d\mu^H.\end{aligned}$$

If in addition the semiflow T_t is mixing, then the formulas in Remark 1.7 hold.

Proposition 2.7 *There is a closed subspace $S \subset C_0^\eta(M, \mathbb{R}^d)$ of infinite codimension such that $\det \Sigma > 0$ for all $v \in C_0^\eta(M, \mathbb{R}^d) \setminus S$.*

Proof Let Σ_Y be as in the proof of Theorem 2.5, so $\Sigma = \bar{H}^{-1} \Sigma_Y$. Suppose that $\det \Sigma = 0$, so $\det \Sigma_Y = 0$. Then there is a nonzero vector $c \in \mathbb{R}^d$ such that $c^T \Sigma_Y c = 0$.

It is standard that $\Sigma_Y = \int_Y m \otimes m d\mu_Y$ (see for example the fourth line of the proof of [18, Theorem 4.3]), so $\int_Y (c \cdot m)^2 d\mu_Y = 0$ and hence $c \cdot m = 0$ a.e. By Proposition 2.3, $c \cdot V = g \circ F - g$ a.e. where $g = c \cdot \chi$. Moreover, it follows from the proof of Proposition 2.3 that $\chi : Y \rightarrow \mathbb{R}^d$ is d_{θ_1} -Lipschitz and hence extends uniquely to a d_{θ_1} -Lipschitz function on Y . It follows that the equation $c \cdot V = g \circ F - g$ extends to the whole of Y .

For $q \geq 1$, let P_q denote the set of points $y \in Y$ such that $F^q y$ is well-defined and $F^q y = y$. Since $F : Y \rightarrow Y$ is full-branch and Markov, such periodic points are dense in Y . Note that $c \cdot \sum_{j=0}^{q-1} V(F^j y) = g(F^q y) - g(y) = 0$ for $y \in P_q$. It follows that $c \cdot \int_0^{H_q(y)} v(T_t y) dt = 0$ where $H_q(y) = \sum_{j=0}^{q-1} H(F^j y)$. Hence, for each periodic orbit $\{y, Fy, \dots, F^{q-1}y\} \subset P_q$, $q \geq 1$, we obtain a distinct nontrivial linear constraint on v . Now define $S = \{v \in C_0^\eta(M, \mathbb{R}^d) : c \cdot \int_0^{H_q(y)} v(T_t y) dt = 0 \text{ for all } y \in P_q, q \geq 1\}$. ■

2.3 Moment estimates

Given $v \in C^\eta(M, \mathbb{R}^d)$, define

$$v_t = \int_0^t v \circ T_s ds, \quad S_t = \int_0^t \int_0^s (v \circ T_r) \otimes (v \circ T_s) dr ds. \quad (2.3)$$

Theorem 2.8 (a) *If $H \in L^p(Y)$ for some $p \geq 2$ and $\inf H > 0$, then there exists $C > 0$ such that*

$$|v_t|_{p-1} \leq Ct^{1/2} \|v\|_\eta \quad \text{for all } v \in C_0^\eta(M, \mathbb{R}^d), t > 0.$$

(b) *If $H \in L^p(Y)$ for some $p \geq 4$ and $\inf H > 0$, then there exists $C > 0$ such that*

$$|S_t|_{(p-1)/2} \leq Ct \|v\|_\eta^2 \quad \text{for all } v \in C_0^\eta(M, \mathbb{R}^d), t > 0.$$

Proof (a) For $t \in [0, 1]$, we have $|v_t|_\infty \leq t|v|_\infty \leq t^{1/2}|v|_\infty$ so we suppose from now on that $t \geq 1$. Recall that $\hat{v} = v \circ \pi : Y^H \rightarrow \mathbb{R}^d$. To prove the result, it is equivalent to show that $|\hat{v}_t|_{L^{p-1}(Y^H)} \leq Ct^{1/2}\|v\|_\eta$ where $\hat{v}_t = \int_0^t \hat{v} \circ F_s ds$.

For $t > 0$, recall from the proof of Theorem 2.5 that the lap number $N(t) : Y^H \rightarrow \mathbb{N}$ is the largest integer $n \geq 0$ such that $u + \sum_{j=0}^{n-1} H(F^j y) \leq t$. Then

$$\hat{v}_t(y, u) = V_{N(t)(y, u)}(y) + \psi \circ F_t(y, u) - \psi(y, u) \quad \text{for } (y, u) \in Y^H,$$

where $\psi(y, u) = \int_0^u v(T_s y) ds$. Also, $|N(t)|_\infty \leq at + 1$ where $a = (\inf H)^{-1} > 0$ by assumption. Let $V_t^*(y) = \sup_{u \in [0, H(y)]} |V_{N(t)(y, u)}(y)|$. By Corollary 2.4, on Y ,

$$|V_t^*|_p \leq \left| \max_{k \leq at+1} |V_k| \right|_p \ll t^{1/2}\|v\|_\eta.$$

Write $p^{-1} + q^{-1} = 1$. By Hölder's inequality,

$$\begin{aligned} \int_{Y^H} |V_{N(t)(y, u)}(y)|^{p-1} d\mu^H(y, u) &\leq \bar{H}^{-1} \int_Y H |V_t^*|^{p-1} d\mu_Y \leq \bar{H}^{-1} |H|_p |V_t^*|_{q(p-1)}^{p-1} \\ &= \bar{H}^{-1} |H|_p |V_t^*|_p^{p-1} \ll t^{(p-1)/2} \|v\|_\eta^{p-1}. \end{aligned}$$

Also, $|\psi \circ F_t|_{p-1} = |\psi|_{p-1}$, and $|\psi| \leq H|v|_\infty$. Using again that $\int_{Y^H} H^{p-1} d\mu^H = \bar{H}^{-1} \int_Y H^p d\mu_Y$, we have that $|\psi \circ F_t|_{p-1} \ll |v|_\infty$. Hence $|\tilde{v}_t|_{p-1} \ll t^{1/2}\|v\|_\eta$.

(b) The steps are analogous to those in part (a). It suffices to show that $|\tilde{S}_t|_{L^{(p-1)/2}(Y^H)} \leq Ct\|v\|_\eta^2$ for $t \geq 1$ where $\tilde{S}_t = \int_0^t \int_0^s (\hat{v} \circ F_r) \otimes (\hat{v} \circ F_s) dr ds$.

First, we establish the corresponding estimate on Y . Let

$$Q_k = \sum_{0 \leq i < j < k} (V \circ F^i) \otimes (V \circ F^j).$$

By Proposition 2.2, $V = m + \chi \circ F - \chi$ where $m \in L^p(Y)$, $\chi \in L^\infty(Y)$ and $m \in \ker P$. Since $V_n = m_n + \chi \circ F^n - \chi$, it is easily checked that

$$\left| \max_{0 \leq k \leq n} |Q_k| \right| = A_1 + A_2,$$

where

$$A_1 = \max_{0 \leq k \leq n} \left| \sum_{0 \leq i < j < k} (m \circ F^i) \otimes (m \circ F^j) \right|, \quad |A_2|_{p/2} \ll n(|m|_{p/2} + |V|_{p/2})|\chi|_\infty \ll n\|v\|_\eta^2.$$

Next,

$$A_1 = \max_{k \leq n} \left| \sum_{i=0}^{k-2} X_{i,k} \right|, \quad X_{i,k} = \left(\sum_{j=1}^{k-i-1} m \otimes (m \circ F^j) \right) \circ F^i.$$

Now $P^{i+1}X_{i,k} = P \sum_{j=1}^{k-i-1} m \otimes (m \circ F^j) = \sum_{j=1}^{k-i-1} (Pm) \otimes (m \circ F^{j-1}) = 0$. This means that $X_{i,k}$ is a sequence of reverse martingale differences with respect to the filtration $\{F^{-i}\mathcal{B}, i \geq 0\}$. Since $p \geq 4$, it follows from Burkholder's inequality that

$$|A_1|_{p/2} \ll n^{1/2} \max_{i < k \leq n} |X_{i,k}|_{p/2}.$$

Applying Burkholder's inequality once more,

$$|X_{i,k}|_{p/2} \leq |m|_p \max_{j \leq n} |m_j|_p \ll n^{1/2} |m|_p^2 \ll n^{1/2} \|v\|_\eta^2.$$

Hence, on Y ,

$$\left| \max_{0 \leq k \leq n} |Q_k| \right|_{p/2} \ll n \|v\|_\eta^2.$$

To pass this estimate from Y to Y^H (and hence M), we use calculations from [18, Proposition 7.5]. Define $\psi(y, u) = \int_0^u v(T_s y) ds$ and the lap number $N(t)$ as in part (a). Then $\hat{v}_t = V_{N(t)} + \psi \circ F_t - \psi$ and

$$\begin{aligned} S_t &= \int_0^t \hat{v}_s \otimes (\hat{v} \circ F_s) ds = \int_0^t V_{N(s)} \otimes (\hat{v} \circ F_s) ds + \int_0^t (\psi \otimes \hat{v}) \circ F_s ds - \psi \otimes \hat{v}_t \\ &= Q_{N(t)} - \hat{v} \otimes \psi + V_{N(t)} \otimes (\psi \circ F_t) + \int_0^t (\psi \otimes \hat{v}) \circ F_s ds - \psi \otimes \hat{v}_t. \end{aligned}$$

Hence working on Y^H , we have

$$\begin{aligned} |S_t|_{(p-1)/2} &\leq |Q_{N(t)}|_{(p-1)/2} + |v|_\infty^2 |H|_{(p-1)/2} + |N(t)|_\infty |V|_{p-1} |v|_\infty |H|_{p-1} \\ &\quad + 2t |v|_\infty^2 |H|_{(p-1)/2} \ll |Q_{N(t)}|_{(p-1)/2} + t |v|_\infty^2. \end{aligned}$$

Finally, setting $Q_t^*(y) = \sup_{u \in [0, H(y)]} |Q_{N(t)(y,u)}(y)|$,

$$\begin{aligned} \int_{Y^H} |Q_{N(t)}|^{(p-1)/2} d\mu^H &\leq \bar{H}^{-1} \int_Y H |Q_t^*|^{(p-1)/2} d\mu_Y \leq \bar{H}^{-1} |H|_{L^p(Y)} |Q_t^*|_{L^{p/2}(Y)}^{(p-1)/2} \\ &\ll \left| \max_{0 \leq k \leq at+1} |Q_k| \right|_{L^{p/2}(Y)}^{(p-1)/2} \ll (t \|v\|_\eta^2)^{(p-1)/2}, \end{aligned}$$

and we obtain $|S_t|_{(p-1)/2} \ll t \|v\|_\eta^2$ as required. \blacksquare

Corollary 2.9 *If $p \geq 4$, then the formulas for Σ and E in Theorem 1.6 hold.*

Proof If $p > 2$ then $W_n(1) = n^{-1/2} v_n \rightarrow W(1)$ by Theorem 2.5 and $|W_n(1)|_q$ is bounded for some $q > 2$ by Theorem 2.8(a). Hence (cf. Corollary 1.5) $\int_M W_n(1) \otimes W_n(1) d\mu_M \rightarrow \mathbb{E}(W(1) \otimes W(1)) = \Sigma$. Similarly, provided that $p \geq 4$, we have $\mathbb{W}_n(1) = n^{-1} S_n \rightarrow_w \mathbb{W}(1)$ and $|\mathbb{W}_n(1)|_q$ bounded for some $q > 1$ so $\int_M \mathbb{W}_n(1) d\mu_M \rightarrow \mathbb{E}(\mathbb{W}(1)) = E$. \blacksquare

3 Rearrangements and improved moment estimates

The results on central limit theorems in Theorem 2.5 are optimal; typically such results fail when H is not L^2 . However, the moment estimates in Theorem 2.8 are far from optimal. To rectify this, we introduce the idea of rearrangement. The strategy is to replace τ by a more unbounded return time in such a way that the flow return time h becomes bounded, while keeping H unchanged.

As in Section 2, $T_t : M \rightarrow M$ is a nonuniformly expanding semiflow with cross-section $X = Y^\tau$ and unbounded roof function $h : X \rightarrow \mathbb{R}^+$. We assume that the induced roof function $H : Y \rightarrow \mathbb{R}^+$ is defined as in (1.1), with $\inf H > 0$, and satisfies (2.1) and (2.2). We assume also that

$$d(T_t x, T_{t'} x) \leq C|t - t'|^\eta \quad \text{for } t, t' \geq 0, x \in M. \quad (3.1)$$

The modified versions $\tilde{\tau}$ and \tilde{h} of τ and h are defined as follows. First,

$$\tilde{\tau} : Y \rightarrow \mathbb{Z}^+, \quad \tilde{\tau}|_{Y_j} = [\|1_{Y_j} H\|_\theta] + 1.$$

Using $F : Y \rightarrow Y$ and $\tilde{\tau} : Y \rightarrow \mathbb{Z}^+$, we construct the tower map

$$f_\Delta : \Delta \rightarrow \Delta, \quad \Delta = Y^{\tilde{\tau}},$$

as in Section 1.2. Now define

$$\tilde{h} : \Delta \rightarrow \mathbb{R}^+, \quad \tilde{h}(y, \ell) = H(y)/\tilde{\tau}(y) \quad \text{for } y \in Y_j, 0 \leq \ell < \tilde{\tau}(y).$$

The key lemma is the following:

Lemma 3.1 (a) $H(y) = \sum_{\ell=0}^{\tilde{\tau}(y)-1} \tilde{h}(y, \ell)$ for all $y \in Y$.

(b) $0 < \inf \tilde{h} \leq |\tilde{h}|_\infty \leq 1$ and

$$|\tilde{h}(y, \ell) - \tilde{h}(y', \ell)| \leq d_\theta(y, y') \quad \text{for } y, y' \in Y_j, j \geq 1, 0 \leq \ell \leq \tilde{\tau}(y) - 1.$$

Proof Part (a) is immediate. Also, $|\tilde{h}(y, \ell) - \tilde{h}(y', \ell)| = |H(y) - H(y')| / \{[\|1_{Y_j} H\|_\theta] + 1\} \leq d_\theta(y, y')$ for $y, y' \in Y_j, j \geq 1, 0 \leq \ell \leq \tilde{\tau}(y) - 1$, and similarly $|\tilde{h}|_\infty \leq 1$.

It remains to verify that \tilde{h} is bounded below. Let $y \in Y_j, 0 \leq \ell < \tilde{\tau}(y)$. Using (2.1) and $\inf H > 0$,

$$\begin{aligned} [\|H1_{Y_j}\|_\theta] + 1 &\leq |1_{Y_j} H|_\infty + |1_{Y_j} H|_\theta + 1 \\ &\leq (2C + 1) \inf_{Y_j} H + (\inf H)^{-1} \inf_{Y_j} H \leq C' \inf_{Y_j} H, \end{aligned}$$

where $C' = 2C + 1 + (\inf H)^{-1}$. Hence $\tilde{h}(y, \ell) = H(y) / ([\|H1_{Y_j}\|_\theta] + 1) \geq 1/C'$ as required. \blacksquare

Since \tilde{h} is bounded above and below, it is immediate from Lemma 3.1(a) that $\mu_Y(H > n) \approx \mu_Y(\tilde{\tau} > n)$. As usual, we assume that H (and $\tilde{\tau}$) is integrable.

As in Section 1.2, we define the suspension semiflow $f_t : \Delta^{\tilde{h}} \rightarrow \Delta^{\tilde{h}}$ and ergodic invariant probability measures $\mu_\Delta = \mu_{\tilde{Y}}$ on Δ and $\mu^{\tilde{h}}$ on $\Delta^{\tilde{h}}$.

Let $\tilde{h}_\ell(y) = \sum_{k=0}^{\ell-1} \tilde{h}(y, k)$ for $\ell \geq 1$ and define

$$\pi_\Delta : \Delta^{\tilde{h}} \rightarrow M, \quad \pi_\Delta(y, \ell, u) = T_{u+\tilde{h}_\ell(y)}y. \quad (3.2)$$

By Lemma 3.1(a), we obtain

Proposition 3.2 π_Δ is a measurable semiconjugacy from $(f_t, \Delta^{\tilde{h}})$ to (T_t, M) .

Proof It suffices to show that π_Δ respects the identifications on Δ and $\Delta^{\tilde{h}}$.

Recall that

$$H(y) = h_{\tau(y)}(y) = \sum_{\ell=0}^{\tau(y)-1} h(f^\ell y), \quad Fy = f^{\tau(y)}y, \quad f(y) = T_{h(y)}y,$$

for $y \in Y$. It follows that

$$F = f^\tau = (T_h)^\tau = T_H \quad \text{on } Y.$$

First, we verify that $\pi_\Delta(y, \tilde{\tau}(y), u) = \pi_\Delta(Fy, 0, u)$. By Lemma 3.1(a),

$$\pi_\Delta(y, \tilde{\tau}(y), u) = T_{u+\tilde{h}_{\tilde{\tau}(y)}(y)}y = T_u T_{H(y)}y = T_u Fy = \pi_\Delta(Fy, 0, u).$$

Second, we verify that $\pi_\Delta(p, \tilde{h}(p)) = \pi_\Delta(f_\Delta p, 0)$. Write $p = (y, \ell) \in \Delta$. If $\ell \leq \tau(y) - 2$, then

$$\begin{aligned} \pi_\Delta(p, \tilde{h}(p)) &= \pi_\Delta(y, \ell, \tilde{h}(y, \ell)) = T_{\tilde{h}(y, \ell)+\tilde{h}_\ell(y)}y \\ &= T_{\tilde{h}_{\ell+1}(y)}y = \pi_\Delta(y, \ell + 1, 0) = \pi_\Delta(f_\Delta(y, \ell), 0). \end{aligned}$$

Also, if $\ell = \tau(y) - 1$, then $\tilde{h}_{\ell+1}(y) = H(y)$ and $f_\Delta(y, \ell) = (Fy, 0)$, so

$$\pi_\Delta(p, \tilde{h}(p)) = T_{H(y)}y = Fy = \pi_\Delta(Fy, 0, 0) = \pi_\Delta(f_\Delta(y, \tau(y) - 1), 0) = \pi_\Delta(f_\Delta p, 0),$$

as required. ■

Remark 3.3 We refer to the procedure in this section as “rearranging the tower” because the obvious suspension flow to consider is $(Y^\tau)^h$ as in Subsection 1.2. As evidenced by Theorem 3.6 below, the rearranged suspension $\Delta^{\tilde{h}} = (Y^{\tilde{\tau}})^{\tilde{h}}$ has significant advantages.

This procedure is analogous to introducing transparent walls, a technique well-known in the theory of billiards (see eg. [22]). When crossing a transparent wall, the velocity of the billiard particle is not modified, nonetheless these events can be regarded as (fake) collisions, which makes the length of the free flight bounded. However, in the present paper, we implement this idea at the more abstract level of the Young tower, which has certain advantages. In particular, there is no need for the fake collision set to be identified with a subset of the ambient phase-space.

3.1 Martingale-coboundary decomposition on Δ

Let $v \in C^\eta(M, \mathbb{R}^d)$ and define $\tilde{v} = v \circ \pi_\Delta : \Delta^{\tilde{h}} \rightarrow \mathbb{R}^d$. Instead of inducing on Y , we induce on Δ , defining the induced observable

$$\tilde{V} : \Delta \rightarrow \mathbb{R}^d, \quad \tilde{V}(y, \ell) = \int_0^{\tilde{h}(y, \ell)} \tilde{v}(y, \ell, u) du.$$

Clearly, $|\tilde{v}|_\infty = |v|_\infty$ and $|\tilde{V}|_\infty \leq |\tilde{h}|_\infty |\tilde{v}|_\infty \leq |v|_\infty$.

Proposition 3.4 *Let $\theta_1 = \theta^{\eta^2}$. There is a constant $C > 0$ such that*

$$|\tilde{V}(y, \ell) - \tilde{V}(y', \ell)| \leq C \|v\|_\eta \tilde{\tau}(y)^\eta d_{\theta_1}(y, y') \quad \text{for } y, y' \in Y_j, j \geq 1, 0 \leq \ell \leq \tilde{\tau}(y) - 1,$$

for all $v \in C^\eta(M, \mathbb{R}^d)$.

Proof Let $(y, \ell, u), (y', \ell, u) \in \Delta^{\tilde{h}}$ with $y, y' \in Y_j, j \geq 1$. Then

$$|\tilde{v}(y, \ell, u) - \tilde{v}(y', \ell, u)| = |v(T_{\tilde{h}_\ell(y)+u} y) - v(T_{\tilde{h}_\ell(y')+u} y')| \leq |v|_\eta d(T_{\tilde{h}_\ell(y)+u} y, T_{\tilde{h}_\ell(y')+u} y')^\eta.$$

By Lemma 3.1(b), $|\tilde{h}_\ell(y) - \tilde{h}_\ell(y')| \leq \sum_{k=0}^{\tilde{\tau}(y)-1} |\tilde{h}(y, k) - \tilde{h}(y', k)| \leq \tilde{\tau}(y) d_\theta(y, y')$, so

$$d(T_{\tilde{h}_\ell(y)+u} y, T_{\tilde{h}_\ell(y')+u} y) \leq C |\tilde{h}_\ell(y) - \tilde{h}_\ell(y')|^\eta \leq C \tilde{\tau}(y)^\eta d_\theta(y, y')^\eta,$$

by (3.1). Also, by (2.2), $d(T_{\tilde{h}_\ell(y')+u} y', T_{\tilde{h}_\ell(y')+u} y') \leq C \theta^{s(y, y')}$. Hence,

$$|\tilde{v}(y, \ell, u) - \tilde{v}(y', \ell, u)| \leq 2C |v|_\eta \tilde{\tau}(y)^\eta d_{\theta_1}(y, y').$$

Next,

$$\begin{aligned} |\tilde{V}(y, \ell) - \tilde{V}(y', \ell)| &\leq \int_{\tilde{h}(y', \ell)}^{\tilde{h}(y, \ell)} |\tilde{v}(y, \ell, u)| du + \int_0^{\tilde{h}(y', \ell)} |\tilde{v}(y, \ell, u) - \tilde{v}(y', \ell, u)| du \\ &\leq |\tilde{h}(y, \ell) - \tilde{h}(y', \ell)| |v|_\infty + |\tilde{h}|_\infty 2C |v|_\eta \tilde{\tau}(y)^\eta d_{\theta_1}(y, y') \ll \|v\|_\eta \tilde{\tau}(y)^\eta d_{\theta_1}(y, y'), \end{aligned}$$

by Lemma 3.1(b). ■

Let $L : L^1(\Delta) \rightarrow L^1(\Delta)$ denote the transfer operator corresponding to $f_\Delta : \Delta \rightarrow \Delta$ (so $\int_\Delta L v w d\mu_\Delta = \int_\Delta v w \circ f_\Delta d\mu_\Delta$ for all $v \in L^1(\Delta), w \in L^\infty(\Delta)$).

Proposition 3.5 *Assume that $\mu_Y(H > n) = O(n^{-\beta})$ for some $\beta > 1$. Let $v \in C_0^\eta(M, \mathbb{R}^d)$. For any $p < \beta - 1$, there exists $\tilde{m}, \tilde{\chi} \in L^p(\Delta, \mathbb{R}^d)$ such that $\tilde{V} = \tilde{m} + \tilde{\chi} \circ f_\Delta - \tilde{\chi}$, $\tilde{m} \in \ker L$.*

Moreover, there exists a constant $C > 0$ such that $|\tilde{m}|_p \leq C \|v\|_\eta$ and $|\tilde{\chi}|_p \leq C \|v\|_\eta$ for all $v \in C_0^\eta(M, \mathbb{R}^d)$.

Proof By Lemma 3.1, $\mu_Y(\tilde{\tau} > n) = O(n^{-\beta})$. Note that

$$\mu_\Delta\{(y, \ell) \in \Delta : \tilde{\tau}(y) > n\} = \left(\int_Y \tilde{\tau} d\mu_Y\right)^{-1} \int_Y \tilde{\tau} 1_{\{\tilde{\tau} > n\}} d\mu_Y = O(n^{-(\beta-1)}). \quad (3.3)$$

Without loss, $d = 1$. Shrinking $\eta \in (0, 1]$ if necessary, we can suppose that $p < \beta - 1 - \eta$.

Fix $n \geq 1$ and let $\tilde{V}_n(y, \ell) = \begin{cases} \tilde{V}(y, \ell) & \tilde{\tau}(y) \leq n \\ 0 & \tilde{\tau}(y) \geq n + 1 \end{cases}$. Then $|\tilde{V}_n|_\infty \leq |\tilde{V}|_\infty \leq |v|_\infty$.

Also, for $y, y' \in Y_j$, $j \geq 1$, it follows from Proposition 3.4 that

$$|\tilde{V}_n(y, \ell) - \tilde{V}_n(y', \ell)| \leq 1_{\{\tilde{\tau}(y) \leq n\}} |\tilde{V}(y, \ell) - \tilde{V}(y', \ell)| \leq C \|v\|_\eta n^\eta d_{\theta_1}(y, y').$$

This means that $\|\tilde{V}_n\| \leq C \|v\|_\eta n^\eta$ in the function space in [37]. Since $\mu_Y(\tilde{\tau} > n) = O(n^{-\beta})$, it follows from [37, Theorem 3] that

$$\left| \int_\Delta \tilde{V}_n w \circ f_\Delta^n d\mu_\Delta - \int_\Delta \tilde{V}_n d\mu_\Delta \int_\Delta w d\mu_\Delta \right| \ll \|\tilde{V}_n\| |w|_\infty n^{-(\beta-1)} \leq C \|v\|_\eta |w|_\infty n^{-(\beta-1-\eta)},$$

for all $w \in L^\infty(\Delta)$, $n \geq 1$. (The dependence on $\|\tilde{V}_n\|$ is not stated explicitly in [37, Theorem 3] but follows by a standard argument using the uniform boundedness principle. Alternatively, see [20] for a direct argument.) Also

$$\left| \int_\Delta (\tilde{V} - \tilde{V}_n) w \circ f_\Delta^n d\mu_\Delta \right| \leq |\tilde{V} - \tilde{V}_n|_\infty |w|_\infty \mu_\Delta\{(y, \ell) : \tilde{\tau}(y) > n\} \ll |v|_\infty |w|_\infty n^{-(\beta-1)},$$

by (3.3), and similarly $\left| \int_\Delta \tilde{V}_n d\mu_\Delta \right| = \left| \int_\Delta (\tilde{V} - \tilde{V}_n) d\mu_\Delta \right| \ll |v|_\infty n^{-(\beta-1)}$ (where we have used that $\int_\Delta \tilde{V} d\mu_\Delta = 0$). Hence

$$\left| \int_\Delta L^n \tilde{V} w d\mu_\Delta \right| = \left| \int_\Delta \tilde{V} w \circ f_\Delta^n d\mu_\Delta \right| \ll \|v\|_\eta |w|_\infty n^{-(\beta-1-\eta)}.$$

Taking $w = \text{sgn } L^n \tilde{V}$, we obtain $|L^n \tilde{V}|_1 \ll \|v\|_\eta n^{-(\beta-1-\eta)}$. Hence

$$|L^n \tilde{V}|_p^p \leq |L^n \tilde{V}|_\infty^{p-1} |L^n \tilde{V}|_1 \leq |\tilde{V}|_\infty^{p-1} |L^n \tilde{V}|_1 \ll |v|_\infty^{p-1} \|v\|_\eta n^{-(\beta-1-\eta)},$$

so $\tilde{\chi} = \sum_{n=1}^\infty L^n \tilde{V}$ converges in L^p . Finally, $L\tilde{m} = L\tilde{V} - \tilde{\chi} + L\tilde{\chi} = 0$. \blacksquare

3.2 Improved moment estimates

Given $v \in C^\eta(M, \mathbb{R}^d)$, define v_t and S_t as in (2.3). A crucial difference from Section 2.3 is that \tilde{V} is L^∞ even though, as before, \tilde{m} is L^p . This enables the use of Rio's inequality [29, 32] following [27]. We continue to assume that $\mu_Y(H > n) = O(n^{-\beta})$ and that $\inf H > 0$.

Theorem 3.6 (a) *Let $2 < p < \beta$. There is a constant $C > 0$ such that*

$$|v_t|_{2(p-1)} \leq Ct^{1/2} \|v\|_\eta \quad \text{for all } v \in C_0^\eta(M, \mathbb{R}^d), t > 0.$$

(b) *Let $4 < p < \beta$. There is a constant $C > 0$ such that*

$$|S_t|_{2(p-1)/3} \leq Ct \|v\|_\eta^2 \quad \text{for all } v \in C_0^\eta(M, \mathbb{R}^d), t > 0.$$

Proof As in the proof of Theorem 2.8, we may suppose that $t \geq 1$. Recall that $\tilde{v} = v \circ \pi_\Delta : \Delta^{\tilde{h}} \rightarrow \mathbb{R}^d$. To prove the result, it is equivalent to show that $|\tilde{v}_t|_{L^{2(p-1)}(\Delta^{\tilde{h}})} \ll t^{1/2} \|v\|_\eta$ and $|\tilde{S}_t|_{L^{2(p-1)/3}(\Delta^{\tilde{h}})} \ll t \|v\|_\eta^2$, where $\tilde{v}_t = \int_0^t \tilde{v} \circ f_s ds$ and $\tilde{S}_t = \int_0^t \int_0^s (\tilde{v} \circ f_r) \otimes (\tilde{v} \circ f_s) dr ds$.

(a) We follow the arguments in [27, Theorem 3.1] and [28, Lemma 4.1]. Recall that $\tilde{V} \in L^\infty(\Delta, \mathbb{R}^d)$. By Proposition 3.5, $\tilde{V} = \tilde{m} + \tilde{\chi} \circ f_\Delta - \tilde{\chi}$ where $\tilde{m}, \tilde{\chi}$ in $L^{p-1}(\Delta, \mathbb{R}^d)$ and $\tilde{m} \in \ker L$. (Unfortunately, $p-1$ here corresponds to p in [27] and Proposition 3.5, but our notation facilitates comparison of Theorems 2.8 and 3.6.) This is the same as the decomposition $\phi = \psi + \chi \circ T - \chi$ in the proof of [27, Theorem 3.1]. Hence it follows just as in the proof of [27, Theorem 3.1] that $|\tilde{V}_n|_{L^{2(p-1)}(\Delta)} \ll n^{1/2} \|v\|_\eta$ where $\tilde{V}_n = \sum_{j=0}^{n-1} \tilde{V} \circ f_\Delta^j$. By [33, Corollary B1] (with $\nu = 2(p-1)$ and $\delta = 1$), we obtain $|\max_{0 \leq k \leq n} |\tilde{V}_k|_{L^{2(p-1)}(\Delta)}| \ll n^{1/2} \|v\|_\eta$.

For $t > 0$, define the lap number $N(t) : \Delta \rightarrow \mathbb{N}$ to be the largest integer $n \geq 0$ such that $\tilde{h}_n \leq t$. By Lemma 3.1(b),

$$|\tilde{v}_t(x, u) - \tilde{V}_{N(t)(x)}(x)| \leq 2|\tilde{v}|_\infty |\tilde{h}|_\infty \leq 2|v|_\infty \quad \text{for } (x, u) \in \Delta^{\tilde{h}}.$$

Also $|N(t)|_\infty \leq at + 1$ where $a = \inf \tilde{h}^{-1} > 0$ by Lemma 3.1(b). Hence

$$\begin{aligned} |\tilde{v}_t|_{L^{2(p-1)}(\Delta^{\tilde{h}})} &\leq 2|v|_\infty + \left(\int_{\Delta^{\tilde{h}}} |\tilde{V}_{N(t)(x)}|^{2(p-1)} d\mu^{\tilde{h}}(x, u) \right)^{1/(2(p-1))} \\ &\leq 2|v|_\infty + \left(\int_\Delta \tilde{h} d\mu_\Delta \right)^{-1/(2(p-1))} |\tilde{h}|_\infty^{1/(2(p-1))} |\tilde{V}_{N(t)}|_{L^{2(p-1)}(\Delta)} \\ &\ll |v|_\infty + \left| \max_{k \leq at+1} |\tilde{V}_k|_{L^{2(p-1)}(\Delta)} \right| \ll |v|_\infty + (at + 1)^{1/2} \|v\|_\eta \ll t^{1/2} \|v\|_\eta. \end{aligned}$$

This completes the proof that $|v_t|_{2(p-1)} \ll t^{1/2} \|v\|_\eta$.

(b) Again, we have $\tilde{V} = \tilde{m} + \tilde{\chi} \circ f_\Delta - \tilde{\chi}$ where $\tilde{V} \in L^\infty(\Delta, \mathbb{R}^d)$, $\tilde{m}, \tilde{\chi}$ in $L^{p-1}(\Delta, \mathbb{R}^d)$ and $\tilde{m} \in \ker L$. This means that the hypotheses of [18, Proposition 7.1] are satisfied and

$$\left| \max_{0 \leq k \leq n} \left| \sum_{i=0}^{k-1} \sum_{j=0}^{i-1} (\tilde{V} \circ f_\Delta^i) \otimes (\tilde{V} \circ f_\Delta^j) \right|_{L^{2(p-1)/3}} \right| \ll n \|v\|_\eta^2.$$

Using that \tilde{h} is bounded above and below, this moment estimate for discrete time again passes over to continuous time, see [18, Proposition 7.5]. \blacksquare

Remark 3.7 The moment estimate in Theorem 3.6(a) is essentially optimal, see [25, 27]. The estimate in Theorem 3.6(b) is the same as the one obtained in [18] and already significantly improves the one in Theorem 2.8(b), but we expect that the optimal estimate is $|S_t|_{p-1} \leq Ct \|v\|_\eta^2$. (Such an estimate is obtained for certain nonuniformly expanding/hyperbolic systems in [21].)

4 Nonuniformly hyperbolic exponentially contracting flows

In this section, we prove our main results for a class of nonuniformly hyperbolic flows. As shown in Section 6, this class of flows includes Lorenz attractors and singular hyperbolic attractors.

4.1 Standing assumptions for flows

Let (M, d) be a bounded metric space with Borel subsets $Y \subset X \subset M$. Suppose that $T_t : M \rightarrow M$, $h : X \rightarrow \mathbb{R}^+$, $f = T_h : X \rightarrow X$ are as in Subsection 1.2. Also, let $\tau : Y \rightarrow \mathbb{Z}^+$ and $F = f^\tau : Y \rightarrow Y$ be as in Subsection 1.2.

We suppose that μ_Y is a Borel ergodic F -invariant probability measure on Y and that τ is integrable on (Y, μ_Y) . Then we obtain an ergodic f -invariant probability measure μ_X on X . We assume that h is integrable on (X, μ_X) and hence obtain an ergodic T_t -invariant probability measure μ_M on M .

Let \mathcal{W}^s be a measurable partition of Y consisting of *stable leaves*. For each $y \in Y$, let W_y^s denote the stable leaf containing y . We require that $F(W_y^s) \subset W_{Fy}^s$ for all $y \in Y$ and that τ is constant on stable leaves.

Let \bar{Y} denote the space obtained from Y after quotienting by \mathcal{W}^s with natural projection $\bar{\pi} : Y \rightarrow \bar{Y}$. We assume that the quotient map $\bar{F} : \bar{Y} \rightarrow \bar{Y}$ is Gibbs-Markov, with partition $\{\bar{Y}_j\}$ and ergodic invariant probability measure $\bar{\mu} = \bar{\pi}_* \mu_Y$. Let $s(y, y')$ denote the separation time on \bar{Y} .

Let $Y_j = \bar{\pi}^{-1} \bar{Y}_j$; these form a partition of Y and each Y_j is a union of stable leaves. The separation time extends to Y , setting $s(y, y') = s(\bar{\pi}y, \bar{\pi}y')$ for $y, y' \in Y$.

4.2 Exponential contraction along stable leaves

Assume the set up in Subsection 4.1. Define the induced roof function $H : Y \rightarrow \mathbb{R}^+$ as in (1.1). We suppose throughout that H is integrable and that $\inf H > 0$.

We require in addition that there is a measurable subset $\hat{Y} \subset Y$ such that for every $y \in Y$ there is a unique $\hat{y} \in \hat{Y} \cap W_y^s$. Assume that there are constants $C \geq 1$, $\gamma_0 \in (0, 1)$, $\eta \in (0, 1]$, such that

$$|H(y) - H(y')| \leq C(\inf_{Y_j} H) \gamma_0^{s(y, y')} \quad \text{for } y, y' \in \hat{Y} \cap Y_j, j \geq 1, \quad (4.1)$$

$$d(T_t y, T_{t'} y) \leq C|t - t'|^\eta \quad \text{for } y \in \hat{Y}, t, t' \geq 0, \quad (4.2)$$

$$d(T_t y, T_{t'} y') \leq C\gamma_0^t \quad \text{for } y \in Y, y' \in W_y^s, t \geq 0, \quad (4.3)$$

$$d(T_t y, T_{t'} y') \leq C\gamma_0^{s(y, y') - J(t)(y)} \quad \text{for } y, y' \in \hat{Y} \cap Y_j, j \geq 1, t \geq 0, \quad (4.4)$$

where $J(t)(y)$ is the largest integer $n \geq 0$ such that $H_n(y) \leq t$ (as usual $H_n = \sum_{j=0}^{n-1} H \circ F^j$). These conditions are standard except for the relatively restrictive

condition (4.3) which represents exponential contraction of stable leaves under the full flow (and not just under the induced map F).

4.3 Martingale-coboundary decomposition

Given $v \in C^n(M, \mathbb{R}^d)$, define the induced observable $V(y) = \int_0^{H(y)} v(T_u y) du$ on Y . Also, define $\chi_1, \widehat{V} : Y \rightarrow \mathbb{R}^d$,

$$\chi_1(y) = \int_0^\infty (v(T_t \hat{y}) - v(T_t y)) dt, \quad \widehat{V} = V + \chi_1 - \chi_1 \circ F.$$

Finally, define $g : \bar{Y} \rightarrow \mathbb{R}$ by $g|_{\bar{Y}_j} = \inf_{Y_j} H$, $j \geq 1$.

Proposition 4.1 *Suppose that $H \in L^p(Y)$ for some $p \geq 1$. There exist constants $C > 0$, $\theta \in (0, 1)$ such that for all $v \in C^n(M, \mathbb{R}^d)$,*

- (a) $\chi_1 \in L^\infty(Y, \mathbb{R}^d)$ and $|\chi_1|_\infty \leq C|v|_\eta$,
- (b) $\widehat{V} = \bar{V} \circ \bar{\pi}$ for some $\bar{V} \in L^p(\bar{Y}, \mathbb{R}^d)$. Moreover, $|\bar{V}| \leq C\|v\|_\eta g$ and $|\bar{V}(y) - \bar{V}(y')| \leq C\|v\|_\eta g(y) d_\theta(y, y')$ for $y, y' \in \bar{Y}_j$, $j \geq 1$.

Proof (a) Let $\gamma_1 = \gamma_0^\eta$. By (4.3),

$$|\chi_1(y)| \leq |v|_\eta \int_0^\infty d(T_t \hat{y}, T_t y)^\eta dt \leq C|v|_\eta \int_0^\infty \gamma_1^t dt \ll |v|_\eta,$$

so $|\chi_1|_\infty \ll |v|_\eta$.

(b) A computation shows that

$$\widehat{V} = V(\hat{y}) + \int_0^\infty (v(T_t F \hat{y}) - v(T_t \widehat{F} y)) dt.$$

In particular, \widehat{V} is constant along stable leaves and has the desired factorisation $\widehat{V} = \bar{V} \circ \bar{\pi}$. It is immediate from the definition of \widehat{V} that $|\widehat{V}| \leq |V| + 2|\chi_1|_\infty \leq |v|_\infty H + 2|\chi_1|_\infty \ll \|v\|_\eta H$. Hence $|\bar{V}| \ll \|v\|_\eta g$ and $\bar{V} \in L^p(\bar{Y}, \mathbb{R}^d)$.

Next, let $y, y' \in Y$ and set $M = s(y, y')/q$ where $q > 1$ is to be specified. Then

$$|\widehat{V}(y) - \widehat{V}(y')| \leq |V(\hat{y}) - V(\hat{y}')| + A_M(y) + A_M(y') + B_M(F\hat{y}, F\hat{y}') + B_M(\widehat{F}y, \widehat{F}y'),$$

where

$$A_M(z) = \int_0^\infty |v(T_t F \hat{z}) - v(T_t \widehat{F} z)| dt, \quad B_M(z, z') = \int_0^M |v(T_t z) - v(T_t z')| dt.$$

The same calculation as in (a) shows that $A_M(z) \ll |v|_\eta \gamma_1^M$ for $z = y, y'$.

Next, let $y, y' \in \widehat{Y} \cap Y_j$. By (4.1), (4.2) and (4.4),

$$\begin{aligned} d(T_t F \hat{y}, T_t F \hat{y}') &= d(T_t F y, T_t F y') = d(T_{t+H(y)} y, T_{t+H(y')} y') \\ &\leq d(T_{t+H(y)} y, T_{t+H(y)} y') + d(T_{t+H(y)} y', T_{t+H(y')} y') \\ &\leq C \gamma_0^{s(y, y') - J(t+H(y))(y)} + |H(y) - H(y')|^\eta \ll (\inf_{Y_j} H) \gamma_1^{s(y, y') - J(t+H(y))(y)}. \end{aligned}$$

Hence, setting $\gamma_2 = \gamma_1^\eta$,

$$B_M(F\hat{y}, F\hat{y}') \leq |v|_\eta \int_0^M d(T_t F\hat{y}, T_t F\hat{y}')^\eta dt \ll |v|_\eta M(\inf_{Y_j} H) \gamma_2^{s(y,y') - J(M+H(y))(y)}.$$

Now $J(t) \leq t/\inf H$ and $J(t + H(y))(y) = 1 + J(t)$. Hence

$$\begin{aligned} B_M(F\hat{y}, F\hat{y}') &\ll |v|_\eta M(\inf_{Y_j} H) \gamma_2^{s(y,y') - (1+M/\inf H)} \\ &\ll |v|_\eta (\inf_{Y_j} H) s(y, y') \gamma_2^{(1-1/(q \inf H))s(y,y')}. \end{aligned}$$

A much simpler calculation shows that

$$B_M(\widehat{Fy}, \widehat{Fy}') \leq C|v|_\eta M \gamma_1^{s(Fy, Fy') - M/\inf H} \ll |v|_\eta s(y, y') \gamma_1^{(1-1/(q \inf H))s(y,y')}.$$

Finally, proceeding as in the proof of Proposition 2.2 and using (4.1) and (4.4),

$$\begin{aligned} |V(\hat{y}) - V(\hat{y}')| &\leq |v|_\infty |H(\hat{y}) - H(\hat{y}')| + C|v|_\eta \min\{H(\hat{y}), H(\hat{y}')\} \gamma_1^{s(y,y')-1} \\ &\ll \|v\|_\eta (\inf_{Y_j} H) \gamma_1^{s(y,y')}. \end{aligned}$$

Altogether, we have shown that $|\widehat{V}(y) - \widehat{V}(y')| \ll \|v\|_\eta (\inf_{Y_j} H) (\gamma_1^{s(y,y')/q} + s(y, y') \gamma_2^{(1-1/(q \inf H))s(y,y')})$. Taking $q = \eta^{-1} + (\inf H)^{-1}$ and $\theta \in (\gamma_1^{1/q}, 1)$, we obtain the desired result. \blacksquare

Corollary 4.2 *Let $v \in C_0^\eta(M, \mathbb{R}^d)$ and suppose that $H \in L^p(Y)$ for some $p \geq 1$. Then there exists $\bar{m} \in L^p(\bar{Y}, \mathbb{R}^d)$, $\chi_1 \in L^\infty(Y, \mathbb{R}^d)$ such that $V = \bar{m} \circ \bar{\pi} + \chi_1 \circ F - \chi_1$ and $\bar{m} \in \ker P$.*

Moreover, there exists $C > 0$ such that $\|\bar{m}\|_p \leq C\|v\|_\eta$ and $\|\chi_1\|_\infty \leq C\|v\|_\eta$ for all $v \in C_0^\eta(M, \mathbb{R}^d)$.

Proof By Proposition 4.1, $V = \widehat{V} + \chi_1 \circ F - \chi_1$ where $\chi_1 \in L^\infty(Y, \mathbb{R}^d)$ and $\widehat{V} = \bar{V} \circ \bar{\pi}$ satisfies the estimates in Proposition 4.1(b). By Proposition 2.1(b), $\|P\bar{V}\|_\theta \ll |g|_1 \leq |H|_1$. Using Proposition 2.1(a), we can proceed as in Proposition 2.3, defining $\chi_2 = \sum_{n=1}^\infty P^n \bar{V} \in L^\infty(\bar{Y}, \mathbb{R}^d)$. Then $\bar{V} = \bar{m} + \chi_2 \circ \bar{F} - \chi_2$ where $\bar{m} \in \ker P$. Now define $\chi = \chi_1 + \chi_2 \circ \bar{\pi}$. \blacksquare

Corollary 4.3 *Suppose that $H \in L^p(Y)$ for some $p \geq 2$. Then there is a constant $C > 0$ such that $\|\max_{1 \leq k \leq n} |V_k|\|_p \leq Cn^{1/2}\|v\|_\eta$ for all $v \in C_0^\eta(M, \mathbb{R}^d)$, $n \geq 1$.*

Proof As in the proof of Corollary 2.4, $\|\max_{1 \leq k \leq n} |(\bar{m} \circ \bar{\pi})_k|\|_{L^p(Y)} = \|\max_{1 \leq k \leq n} |\bar{m}_k|\|_{L^p(\bar{Y})} \ll n^{1/2}$. The result follows. \blacksquare

4.4 Central limit theorems and moment estimates

We continue to assume the set up in Subsections 4.1 and 4.2, including that $\inf H > 0$.

Theorem 4.4 *Suppose that $H \in L^2(Y)$. Let $v \in C_0^\eta(M, \mathbb{R}^d)$. Then the statement and conclusion of Theorem 2.5 holds. That is, $(W_n, \mathbb{W}_n) \rightarrow_w (W, \mathbb{W})$ in $C([0, \infty), \mathbb{R}^d \times \mathbb{R}^{d \times d})$ on (M, μ_M) , where W_n, \mathbb{W}_n, W and \mathbb{W} are defined as in Theorem 2.5.*

Proof Define W_n^Y, \mathbb{W}_n^Y on (Y, μ_Y) as in the proof of Theorem 2.5. In place of Proposition 2.3 and [18, Theorem 4.3], we use Corollary 4.2 and [18, Theorem 5.2] to deduce that $(W_n^Y, \mathbb{W}_n^Y) \rightarrow (W^Y, \mathbb{W}^Y)$ in $D([0, \infty), \mathbb{R}^d \times \mathbb{R}^{d \times d})$ where W^Y and \mathbb{W}^Y are as in step (i) of the proof of Theorem 2.5.

Claims (a)–(c) are verified in exactly the same way as in the proof of Theorem 2.5 (with Corollary 2.4 replaced by Corollary 4.3). The remainder of the proof of step (ii) is identical to Theorem 2.5, again using [18, Theorem 6.1] to pass the iterated WIP from Y to Y^H and hence M . ■

Theorem 4.5 *Define v_t, S_t as in (2.3).*

- (a) *If $H \in L^p(Y)$ for some $p \geq 2$, then there is a constant $C > 0$ such that $|v_t|_{p-1} \leq Ct^{1/2} \|v\|_\eta$ for all $v \in C_0^\eta(M, \mathbb{R}^d)$, $t > 0$.*
- (b) *If $H \in L^p(Y)$ for some $p \geq 4$, then there is a constant $C > 0$ such that $|S_t|_{(p-1)/2} \leq Ct \|v\|_\eta^2$ for all $v \in C_0^\eta(M, \mathbb{R}^d)$, $t > 0$.*

Proof The proof is identical to that of Theorem 2.8. ■

Remark 4.6 (i) Again, it follows that the formulas for Σ and E in Theorem 1.6 hold provided $p \geq 4$, and that $\det \Sigma > 0$ for v lying outside a closed subspace of infinite codimension.

(ii) Under appropriate hypotheses on the flow, it is possible to incorporate the rearrangement idea from Section 3 and hence to obtain the improved moment estimates in Theorem 3.6. We omit the tedious details.

5 Nonuniformly hyperbolic flows with horseshoe structure

In this section, we relax the assumption (4.3) about exponential contraction along stable leaves under the flow. Instead we assume that $F : Y \rightarrow Y$ is a nonuniformly hyperbolic horseshoe as in [36]. We begin by describing the basic set up, focusing on the parts needed for understanding what follows and referring to [36] for further details.

We assume the structure from Section 4.1, including the measurable partition \mathcal{W}^s of Y into stable leaves. Further, we suppose that there is a second measurable partition \mathcal{W}^u of Y consisting of *unstable leaves*.

We assume that there are constants $C > 0$, $\gamma_0 \in (0, 1)$, $\eta \in (0, 1]$, such that for all $y \in Y$, $y' \in W_y^s$, $t \in [0, H(y)] \cap [0, H(y')]$,

$$|H(y) - H(y')| \leq C(\inf_{Y_j} H)d(y, y')^\eta, \quad d(T_t y, T_t y') \leq C d(y, y')^\eta; \quad (5.1)$$

and for all $y, y' \in Y_j$, $j \geq 1$, $y' \in W_y^u$, $t \in [0, H(y)] \cap [0, H(y')]$,

$$|H(y) - H(y')| \leq C(\inf_{Y_j} H)\gamma_0^{s(y, y')}, \quad d(T_t y, T_t y') \leq C\gamma_0^{s(y, y')}. \quad (5.2)$$

Also, we assume that for all $y, y' \in Y_j$, $j \geq 1$, $n \geq 0$,

$$d(F^n y, F^n y') \leq C\gamma_0^n \quad \text{for } y' \in W_y^s, \quad (5.3)$$

$$d(y, y') \leq C\gamma_0^{s(y, y')} \quad \text{for } y' \in W_y^u. \quad (5.4)$$

Define the induced observable $V(y) = \int_0^{H(y)} v(T_u y) du$ on Y .

Proposition 5.1 *Let $v \in C^\eta(M, \mathbb{R}^d)$, and set $\gamma_1 = \gamma_0^\eta$. Then $|V| \leq |v|_\infty H$ and there is a constant $C > 0$ such that $|V(y) - V(y')| \leq C\|v\|_\eta(\inf_{Y_j} H)d(y, y')^{\eta^2}$ for $y \in Y$, $y' \in W_y^s$; while $|V(y) - V(y')| \leq C\|v\|_\eta(\inf_{Y_j} H)\gamma_1^{s(y, y')}$ for $y, y' \in Y_j$, $j \geq 1$, $y' \in W_y^u$.*

Proof The first estimate is immediate. Let $y' \in W_y^s$ and suppose without loss that $H(y) \geq H(y')$. Then

$$\begin{aligned} |V(y) - V(y')| &\leq \int_{H(y')}^{H(y)} |v(T_u y)| du + \int_0^{H(y')} |v(T_u y) - v(T_u y')| du \\ &\leq |v|_\infty(H(y) - H(y')) + |v|_\eta H(y') \max_{u \in [0, H(y')]} d(T_u y, T_u y')^\eta \ll \|v\|_\eta (\inf_{Y_j} H) d(y, y')^{\eta^2}, \end{aligned}$$

by (5.1). The last estimate follows by the same argument using (5.2) instead of (5.1). \blacksquare

5.1 Martingale-coboundary decomposition

We apply a Gordin type argument [15] to obtain an L^2 martingale-coboundary decomposition. Let $\bar{\mathcal{B}}$ denote the underlying σ -algebra on \bar{Y} and let $\mathcal{B} = \bar{\pi}^{-1}\bar{\mathcal{B}}$. By Proposition 5.1, $|\mathbb{E}(V|\mathcal{B})|_1 \leq |V|_1 \leq |v|_\infty |H|_1$. Since $\mathbb{E}(V|\mathcal{B})$ is constant along stable leaves, we can write $\mathbb{E}(V|\mathcal{B}) = \bar{V} \circ \bar{\pi}$ where $\bar{V} \in L^1(\bar{Y})$.

Let P be the transfer operator for the Gibbs-Markov map $\bar{F} : \bar{Y} \rightarrow \bar{Y}$.

Proposition 5.2 *Suppose that $H \in L^p(Y)$ for some $p \geq 1$. There is a constant $C > 0$ such that*

$$\sum_{n \geq 0} |\mathbb{E}(V \circ F^n | \mathcal{B}) - V \circ F^n|_{L^p(Y)} \leq C \|v\|_\eta, \quad \sum_{n \geq 1} |P^n \bar{V}|_{L^\infty(\bar{Y})} \leq C \|v\|_\eta,$$

for all $v \in C_0^\eta(M, \mathbb{R}^d)$.

Proof Set $\gamma_2 = \gamma_0^2$. Let $y, y' \in Y$ with $y' \in W_y^s$. Then

$$|V(F^n y) - V(F^n y')| \ll \|v\|_\eta H(F^n y) d(F^n y, F^n y')^{\eta^2} \ll \gamma_2^n \|v\|_\eta H(F^n y),$$

by Proposition 5.1 and (5.3). Hence $|\mathbb{E}(V \circ F^n | \mathcal{B})(y) - V(F^n y)| \ll \gamma_2^n \|v\|_\eta H(F^n y)$ and

$$|\mathbb{E}(V \circ F^n | \mathcal{B}) - V \circ F^n|_p \ll \gamma_2^n \|v\|_\eta \|H\|_p,$$

so the first sum converges.

Now define $g \in L^1(Y)$ and $\bar{g} \in L^1(\bar{Y})$ by setting $g(y) = \inf_{Y_j} H$ for all $y \in Y_j$ and writing $g = \bar{g} \circ \bar{\pi}$. By Proposition 5.1, we have $|V| \ll |v|_\infty g$ and $|V(y) - V(y')| \ll \|v\|_\eta g(y) \gamma_1^{s(y, y')}$ for all $y, y' \in Y_j$, $j \geq 1$, with $y' \in W_y^u$. It follows that $|\bar{V}| \ll |v|_\infty \bar{g}$. Arguing as in [36, Sublemma, p. 612] (see [11, Lemma 5.4]), $|\bar{V}(y) - \bar{V}(y')| \ll \|v\|_\eta \bar{g}(y) \gamma_1^{s(y, y')}$ for $y, y' \in \bar{Y}_j$ (cf. [11, Lemma 5.4]). By Proposition 2.1, there exists $\gamma \in (0, 1)$ such that $|P^n \bar{V}|_\infty \ll \gamma^n \|P \bar{V}\|_{\gamma_1} \ll \gamma^n \|v\|_\eta$ for $n \geq 1$. Hence $\sum_{n \geq 1} |P^n \bar{V}|_\infty < \infty$, completing the proof. \blacksquare

Corollary 5.3 *Let $v \in C_0^\eta(M, \mathbb{R}^d)$ and suppose that $H \in L^p(Y)$ for some $p \geq 1$. There exists $\bar{m} \in L^p(\bar{Y}, \mathbb{R}^d)$, $\chi \in L^p(Y, \mathbb{R}^d)$ such that $V = \bar{m} \circ \bar{\pi} + \chi \circ F - \chi$ and $\bar{m} \in \ker P$.*

Moreover, there is a constant $C > 0$ such that $|\bar{m}|_p \leq C \|v\|_\eta$ and $|\chi|_p \leq C \|v\|_\eta$ for all $v \in C_0^\eta(M, \mathbb{R}^d)$.

Proof Define $g_n = \mathbb{E}(V \circ F^n | \mathcal{B})$ and $h_n = (P^n \bar{V}) \circ \bar{\pi}$. By Proposition 5.2,

$$\chi = \sum_{n \geq 0} (g_n - V \circ F^n) + \sum_{n \geq 1} h_n$$

converges in $L^p(Y)$. Define $m = V + \chi - \chi \circ F \in L^p(Y)$. We claim that (i) m is \mathcal{B} -measurable and (ii) $\mathbb{E}(m | F^{-1} \mathcal{B}) = 0$. Property (i) ensures that $m = \bar{m} \circ \bar{\pi}$ where $\bar{m} : \bar{Y} \rightarrow \mathbb{R}^d$. By property (ii),

$$(UP\bar{m}) \circ \bar{\pi} = \mathbb{E}(\bar{m} | \bar{F}^{-1} \bar{\mathcal{B}}) \circ \bar{\pi} = \mathbb{E}(\bar{m} \circ \bar{\pi} | \bar{\pi}^{-1} \bar{F}^{-1} \bar{\mathcal{B}}) = \mathbb{E}(m | F^{-1} \mathcal{B}) = 0,$$

where $Uv = v \circ \bar{F}$. Hence $UP\bar{m} = 0$ and so $\bar{m} \in \ker P$.

It remains to verify the claim. We have

$$m = \sum_{n \geq 0} g_n - \sum_{n \geq 0} g_n \circ F + \sum_{n \geq 1} h_n - \sum_{n \geq 1} h_n \circ F. \quad (5.5)$$

Clearly, $g_n = \mathbb{E}(V \circ F^n | \mathcal{B})$ is \mathcal{B} -measurable. Also, $P^n \bar{V}$ is $\bar{\mathcal{B}}$ -measurable, so $h_n = (P^n \bar{V}) \circ \bar{\pi}$ is measurable with respect to $\bar{\pi}^{-1} \bar{\mathcal{B}} = \mathcal{B}$. Next, $g_n \circ F$ is measurable with respect to $F^{-1} \mathcal{B} \subset \mathcal{B}$ and similarly for $h_n \circ F$. Part (i) of the claim follows.

For part (ii) of the claim, note that $g_n \circ F = \mathbb{E}(V \circ F^n | \mathcal{B}) \circ F = \mathbb{E}(V \circ F^{n+1} | F^{-1} \mathcal{B})$. Hence

$$\mathbb{E}(g_n \circ F | F^{-1} \mathcal{B}) = \mathbb{E}(V \circ F^{n+1} | F^{-1} \mathcal{B}) = \mathbb{E}(g_{n+1} | F^{-1} \mathcal{B}).$$

Also,

$$h_n \circ F = (P^n \bar{V}) \circ \bar{F} \circ \bar{\pi} = (UP^n \bar{V}) \circ \bar{\pi} = \mathbb{E}(P^{n-1} \bar{V} | \bar{F}^{-1} \bar{\mathcal{B}}) \circ \bar{\pi} = \mathbb{E}(h_{n-1} | F^{-1} \mathcal{B}),$$

so $\mathbb{E}(h_n \circ F | F^{-1} \mathcal{B}) = \mathbb{E}(h_{n-1} | F^{-1} \mathcal{B})$. Substituting into (5.5), $\mathbb{E}(m | F^{-1} \mathcal{B}) = \mathbb{E}(g_0 | F^{-1} \mathcal{B}) - \mathbb{E}(h_0 | F^{-1} \mathcal{B})$. Finally, $h_0 = \bar{V} \circ \bar{\pi} = \mathbb{E}(V | \mathcal{B}) = g_0$, so $\mathbb{E}(m | F^{-1} \mathcal{B}) = 0$ as required. \blacksquare

Corollary 5.4 *Suppose that $H \in L^p(Y)$ for some $p \geq 2$. There exists $C > 0$ such that $\left| \max_{1 \leq k \leq n} |V_k| \right|_p \leq Cn^{1/2} \|v\|_\eta$ for all $v \in C_0^\eta(M, \mathbb{R}^d)$, $n \geq 1$.*

Proof As in the proof of Corollary 2.4, $\left| \max_{1 \leq k \leq n} |(\bar{m} \circ \bar{\pi})_k| \right|_{L^p(Y)} = \left| \max_{1 \leq k \leq n} |\bar{m}_k| \right|_{L^p(\bar{Y})} \ll n^{1/2} |\bar{m}|_{L^p(\bar{Y})} \ll n^{1/2} \|v\|_\eta$. Also, $\int_Y \max_{k \leq n} |\chi|^p \circ F^k d\mu_Y \leq \int_Y \sum_{k \leq n} |\chi|^p \circ F^k d\mu_Y = n |\chi|_p^p \ll n \|v\|_\eta^p$. Since $p \geq 2$, the result follows. \blacksquare

5.2 Central limit theorems and moment estimates

We continue to assume the set up from the beginning of the section.

Theorem 5.5 *Suppose that $H \in L^2(Y)$. Let $v \in C_0^\eta(M, \mathbb{R}^d)$. Then the statement and conclusion of Theorem 2.5 holds. That is, $(W_n, \mathbb{W}_n) \rightarrow_w (W, \mathbb{W})$ in $C([0, \infty), \mathbb{R}^d \times \mathbb{R}^{d \times d})$ on (M, μ_M) , where W_n, \mathbb{W}_n, W and \mathbb{W} are defined as in Theorem 2.5.*

Proof Define W_n^Y, \mathbb{W}_n^Y on (Y, μ_Y) as in the proof of Theorem 2.5. In place of Proposition 2.3 and [18, Theorem 4.3], we use Corollary 5.3 and [18, Theorem 5.2] to deduce that $(W_n^Y, \mathbb{W}_n^Y) \rightarrow (W^Y, \mathbb{W}^Y)$ in $D([0, \infty), \mathbb{R}^d \times \mathbb{R}^{d \times d})$ where W^Y and \mathbb{W}^Y are as in step (i) of the proof of Theorem 2.5.

Claims (a)–(c) are verified in exactly the same way as in the proof of Theorem 2.5 (with Corollary 2.4 replaced by Corollary 5.4). The remainder of the proof of step (ii) is identical to Theorem 2.5, again using [18, Theorem 6.1] to pass the iterated WIP from Y to Y^H and hence M . \blacksquare

Theorem 5.6 *Define v_t, S_t as in (2.3). Suppose that $\inf H > 0$.*

(a) *If $H \in L^p(Y)$ for some $p \geq 2$, then there is a constant $C > 0$ such that $|v_t|_{p-1} \leq Ct^{1/2} \|v\|_\eta$ for all $v \in C_0^\eta(M, \mathbb{R}^d)$, $t > 0$.*

(b) If $H \in L^p(Y)$ for some $p \geq 4$, then there is a constant $C > 0$ such that $|S_t|_{(p-1)/2} \leq Ct \|v\|_\eta^2$ for all $v \in C_0^\eta(M, \mathbb{R}^d)$, $t > 0$.

Proof The proof is almost identical to that of Theorem 2.8. The only difference is that $\chi \in L^p(Y, \mathbb{R}^d)$ (instead of $\chi \in L^\infty(Y, \mathbb{R}^d)$) but it is easily checked that the required estimates go through. \blacksquare

6 Singular hyperbolic attractors

In this section, we verify that our main results apply to codimension two singular hyperbolic attractors, as described in the introduction.

Lemma 6.1 *Codimension two singular hyperbolic attractors (Definition 1.2) satisfy the hypotheses in Section 4 with β arbitrarily large.*

Proof We apply the construction in [8], representing the singular hyperbolic attractor Λ as a suspension $\Lambda = X^h$, $X = Y^\tau$, where $\mu_X(h > t) = O(e^{-ct})$ for some $c > 0$ and $\mu_Y(\tau > n) = O(n^{-\beta})$ with β arbitrarily large. By Remark A.2(c), the induced roof function $H = h_\tau$ satisfies $\mu_Y(H > t) = O(t^{-\beta})$ with β arbitrarily large. (In [8], our h and τ are denoted τ and ρ respectively.)

Singular hyperbolic attractors have the property that there is a Hölder stable foliation for the flow [8, Theorem 6.2] and hence the projection $\bar{\pi} : Y \rightarrow \bar{Y}$ extends to $\bar{\pi} : X \rightarrow \bar{X}$. As in [8, Section 9], we can suppose without loss that h and H are constant along stable leaves for the flow and hence are well-defined on \bar{X} and \bar{Y} respectively. Define the quotient Poincaré map $\bar{f} : \bar{X} \rightarrow \bar{X}$. By [8, Lemma 8.4], the induced quotient map $\bar{F} = \bar{f}^\tau : \bar{Y} \rightarrow \bar{Y}$ is Gibbs-Markov.

By [8, Proposition 8.6],

$$|H(y) - H(y')| \leq \sum_{\ell=0}^{\tau(y)-1} |h(\bar{f}^\ell y) - h(\bar{f}^\ell y')| \ll |\bar{F}y - \bar{F}y'|^\epsilon \quad \text{for } y, y' \in \bar{Y}_j, j \geq 1.$$

Hence $|H(y) - H(y')| \ll \gamma_0^{s(y,y')}$ for $y, y' \in \bar{Y}_j, j \geq 1$. It follows immediately that $|H(y) - H(y')| \ll \gamma_0^{s(y,y')}$ for $y, y' \in Y_j, j \geq 1$, so condition (4.1) is satisfied. Also, since the stable foliation is Hölder,

$$|H(y) - H(y')| \ll |Fy - Fy'|^{\epsilon^2} \quad \text{for } y, y' \in Y_j, j \geq 1. \quad (6.1)$$

Condition (4.2) is immediate from smoothness of the flow (with $\eta = 1$). Singular hyperbolic flows contract exponentially along stable leaves by definition, so (4.3) holds. Finally we verify (4.4). There are three steps.

Step 1. We show that $|T_u x - T_u x'| \ll |x - x'|^\epsilon + |fx - fx'|^\epsilon$ for all $x, x' \in X$, $u \in [0, h(x)]$. The first return time h is either bounded or of the form $h = h_1 + h_2$ where h_1 is bounded and h_2 is the first hit map in a neighbourhood of a steady-state

for the flow. The claim is clear for bounded return times, so it remains to consider the first hit map h_2 . Again, the claim is immediate if the flow near the equilibrium is linear. By a Hölder version of the Hartman-Grobman theorem, linearity always holds up to a bi-Hölder change of coordinates completing the proof of Step 1.

Step 2. We show that $|T_t y - T_t y'| \ll |Fy - Fy'|^{\epsilon^2}$ for all $y, y' \in \widehat{Y}_j$, $j \geq 1$, $t \in [0, H(y)] \cap [0, H(y')]$. There exists $0 \leq \ell, \ell' \leq \tau(y) - 1$ such that $t \in [h_\ell(y), h_{\ell+1}(y)] \cap [h_{\ell'}(y'), h_{\ell'+1}(y')]$. Without loss $\ell \leq \ell'$. Then $t = h_\ell(y) + u = h_\ell(y') + u'$ where $u \in [0, h(f^\ell(y))]$ and $u' = t - h_\ell(y') \geq t - h_{\ell'}(y') \geq 0$. By Step 1 and smoothness of the flow,

$$\begin{aligned} |T_t y - T_t y'| &= |T_u f^\ell y - T_{u'} f^\ell y'| \leq |T_u f^\ell y - T_u f^\ell y'| + |T_u f^\ell y' - T_{u'} f^\ell y'| \\ &\ll |f^\ell y - f^\ell y'|^\epsilon + |f^{\ell+1} y - f^{\ell+1} y'|^\epsilon + |u - u'| \\ &= |f^\ell y - f^\ell y'|^\epsilon + |f^{\ell+1} y - f^{\ell+1} y'|^\epsilon + |h_\ell(y) - h_\ell(y')|. \end{aligned}$$

By (6.1), $|h_\ell(y) - h_\ell(y')| \ll |Fy - Fy'|^{\epsilon^2}$. By the construction in [8] which is based on [4, 5], $|f^\ell y - f^\ell y'| \ll |\bar{F}y - \bar{F}y'|^\epsilon \ll |Fy - Fy'|^{\epsilon^2}$ for $\ell \leq \tau(y)$. Hence, $|T_t y - T_t y'| \ll |Fy - Fy'|^{\epsilon^2}$ as required.

Step 3. We complete the verification of (4.4). Let $y, y' \in \widehat{Y} \cap Y_j$. Let $n = J(t)(y)$, $n' = J(t)(y')$ where $J(t)$ is as in (4.4), and suppose without loss that $n \leq n'$ and $n \leq s(y, y')$. Write $T_t y = T_u F^n y$, $T_t y' = T_{u'} F^n y'$ where $u \in [0, H(F^n y)]$ and $u' \geq 0$. By Step 2 and existence of an unstable cone field [8, Proposition 4.2], $|T_u F^n y - T_{u'} F^n y'| \ll |F^{n+1} y - F^{n+1} y'|^\epsilon$. Using this and smoothness of the flow,

$$\begin{aligned} |T_t y - T_t y'| &= |T_u F^n y - T_{u'} F^n y'| \leq |T_u F^n y - T_u F^n y'| + |T_u F^n y' - T_{u'} F^n y'| \\ &\ll |F^{n+1} y - F^{n+1} y'|^\epsilon + |H_n(y) - H_n(y')|. \end{aligned}$$

By (6.1), $|T_t y - T_t y'| \ll \sum_{j=1}^{n+1} |F^j y - F^j y'|^{\epsilon^2}$. By [8, Proposition 4.2], there exists $\lambda_1 \in (0, 1)$ such that $|\bar{F}^{s(y, y')} y - \bar{F}^{s(y, y')} y'| \geq \lambda_1^{-(s(y, y')-j)} |F^j y - F^j y'|$, so $|T_t y - T_t y'| \ll \gamma_0^{s(y, y')-n} = \gamma_0^{s(y, y')-J(t)(y)}$ as required with $\gamma_0 = \lambda_1^{\epsilon^2}$. ■

A Tails for induced roof functions

For completeness, we state and prove the following standard result.

Proposition A.1 $\mu_Y(H > t) \leq \mu_Y(\tau > n) + \bar{\tau} \mu_X(h > t/n)$ for all $n \geq 1$, $t > 1$.

Proof Recall that $\mu^\tau = c^{-1}(\mu_Y \times \text{counting})$, $c = \bar{\tau}$. Also, $\mu_X = \pi_* \mu^\tau$ where $\pi(y, \ell) = f^\ell y$.

Let $h^*(y) = \max\{h(y), h(fy), \dots, h(f^{\tau(y)-1}y)\}$ and define

$$\ell(y) = \min\{\ell \in \{0, \dots, \tau(y) - 1\} : h(f^\ell y) = h^*(y)\}.$$

Then

$$\begin{aligned}
\mu_Y\{H > t, \tau \leq n\} &\leq \mu_Y(h^* > t/n) = \mu_Y\{y \in Y : h(f^{\ell(y)}y) > t/n\} \\
&= c\mu^\tau\{(y, 0) \in Y^\tau : h(f^{\ell(y)}y) > t/n\} = c\mu^\tau\{(y, \ell(y)) \in Y^\tau : h(f^{\ell(y)}y) > t/n\} \\
&= c\mu^\tau\{(y, \ell(y)) \in Y^\tau : h \circ \pi(y, \ell(y)) > t/n\} \leq c\mu^\tau\{p \in Y^\tau : h \circ \pi(p) > t/n\} \\
&= c\mu_X\{x \in X : h(x) > t/n\}.
\end{aligned}$$

Hence $\mu_Y(H > t) \leq \mu_Y(\tau > n) + \mu_Y(H > t, \tau \leq n) \leq \mu_Y(\tau > n) + c\mu_X(h > t/n)$ as required. \blacksquare

Remark A.2 For completeness, we list some special cases of Proposition A.1:

- (a) If $\mu_X(h > t) = O(t^{-\beta})$ and $\mu_Y(\tau > n)$ decays exponentially, then $\mu_Y(H > t) = O((\log t)^\beta t^{-\beta})$. (Choose $n = C \log t$ with C large.)
- (b) If $\mu_X(h > t) = O(e^{-ct})$ and $\mu_Y(\tau > n) = O(e^{-dn})$, then $\mu_Y(H > t) = O(e^{-(cdt)^{\frac{1}{2}}})$. (Choose $n = \sqrt{ct/d}$.)
- (c) If $\mu_X(h > t) = O(e^{-ct})$ and $\mu_Y(\tau > n) = O(n^{-\beta})$, then $\mu_Y(H > t) = O((\log t)^\beta t^{-\beta})$. (Choose $n = at(\log t)^{-1}$ with a small.)
- (d) If $\mu_X(h > t) = O(t^{-\beta_1})$ and $\mu_Y(\tau > n) = O(n^{-\beta_2})$, then $\mu_Y(H > t) = O(t^{-\beta})$ where $\beta = \frac{\beta_1\beta_2}{\beta_1 + \beta_2}$. (Choose $n = t^{\beta_1/(\beta_1+\beta_2)}$.)

Acknowledgements The research of PB was supported in part by Hungarian National Foundation for Scientific Research (NKFIH OTKA) grants K104745 and K123782. The research of IM was supported in part by a European Advanced Grant *StochExtHomog* (ERC AdG 320977). We are grateful to the referee for very helpful suggestions.

References

- [1] J. Aaronson. *An Introduction to Infinite Ergodic Theory*. Math. Surveys and Monographs **50**, Amer. Math. Soc., 1997.
- [2] J. Aaronson and M. Denker. Local limit theorems for partial sums of stationary sequences generated by Gibbs-Markov maps. *Stoch. Dyn.* **1** (2001) 193–237.
- [3] V. S. Afraïmovič, V. V. Bykov and L. P. Silnikov. The origin and structure of the Lorenz attractor. *Dokl. Akad. Nauk SSSR* **234** (1977) 336–339.
- [4] J. F. Alves, J. M. Freitas, S. Luzzatto and S. Vaienti. From rates of mixing to recurrence times via large deviations. *Adv. Math.* **228** (2011) 1203–1236.

- [5] J. F. Alves, S. Luzzatto and V. Pinheiro. Markov structures and decay of correlations for non-uniformly expanding dynamical systems. *Ann. Inst. H. Poincaré Anal. Non Linéaire* **22** (2005) 817–839.
- [6] V. Araújo and I. Melbourne. Exponential decay of correlations for nonuniformly hyperbolic flows with a $C^{1+\alpha}$ stable foliation, including the classical Lorenz attractor. *Ann. Henri Poincaré* **17** (2016) 2975–3004.
- [7] V. Araújo and I. Melbourne. Existence and smoothness of the stable foliation for sectional hyperbolic attractors. *Bull. London Math. Soc.* **49** (2017) 351–367.
- [8] V. Araújo and I. Melbourne. Mixing properties and statistical limit theorems for singular hyperbolic flows without a smooth stable foliation. Preprint, 2017.
- [9] V. Araujo, M. J. Pacifico, E. R. Pujals and M. Viana. Singular-hyperbolic attractors are chaotic. *Trans. Amer. Math. Soc.* **361** (2009) 2431–2485.
- [10] P. Bálint, O. Butterley and I. Melbourne. Polynomial decay of correlations for flows, including Lorentz gas examples. Preprint, 2017.
- [11] P. Bálint and S. Gouëzel. Limit theorems in the stadium billiard. *Comm. Math. Phys.* **263** (2006) 461–512.
- [12] M. Benedicks and L. Carleson. On iterations of $1 - ax^2$ on $(-1, 1)$. *Ann. of Math.* **122** (1985) 1–25.
- [13] D. L. Burkholder. Distribution function inequalities for martingales. *Ann. Probability* **1** (1973) 19–42.
- [14] F. Dumortier, H. Kokubu and H. Oka. A degenerate singularity generating geometric Lorenz attractors. *Ergodic Theory Dynam. Systems* **15** (1995) 833–856.
- [15] M. I. Gordin. The central limit theorem for stationary processes. *Soviet Math. Dokl.* **10** (1969) 1174–1176.
- [16] J. Guckenheimer and R. F. Williams. Structural stability of Lorenz attractors. *Inst. Hautes Études Sci. Publ. Math.* **50** (1979) 59–72.
- [17] M. Holland and I. Melbourne. Central limit theorems and invariance principles for Lorenz attractors. *J. London Math. Soc.* **76** (2007) 345–364.
- [18] D. Kelly and I. Melbourne. Smooth approximation of stochastic differential equations. *Ann. Probab.* **44** (2016) 479–520.
- [19] D. Kelly and I. Melbourne. Homogenization for deterministic fast-slow systems with multidimensional multiplicative noise. *J. Funct. Anal.* **272** (2017) 4063–4102.

- [20] A. Korepanov, Z. Kosloff and I. Melbourne. Explicit coupling argument for nonuniformly hyperbolic transformations. *Proc. Roy. Soc. Edinburgh A*. To appear.
- [21] A. Korepanov, Z. Kosloff and I. Melbourne. Deterministic homogenization for families of fast-slow systems. In preparation.
- [22] A. Krámli, N. Simányi and D. Szász. A “transversal” fundamental theorem for semi-dispersing billiards. *Comm. Math. Phys.* **129** (1990) 535–560.
- [23] R. Leplaideur and D. Yang. SRB measures for higher dimensional singular partially hyperbolic attractors. *Ann. Inst. Fourier (Grenoble)* **67** (2017) 2703–2717.
- [24] E. D. Lorenz. Deterministic nonperiodic flow. *J. Atmosph. Sci.* **20** (1963) 130–141.
- [25] I. Melbourne. Large and moderate deviations for slowly mixing dynamical systems. *Proc. Amer. Math. Soc.* **137** (2009) 1735–1741.
- [26] I. Melbourne and M. Nicol. Almost sure invariance principle for nonuniformly hyperbolic systems. *Comm. Math. Phys.* **260** (2005) 131–146.
- [27] I. Melbourne and M. Nicol. Large deviations for nonuniformly hyperbolic systems. *Trans. Amer. Math. Soc.* **360** (2008) 6661–6676.
- [28] I. Melbourne and A. Török. Convergence of moments for Axiom A and nonuniformly hyperbolic flows. *Ergodic Theory Dynam. Systems* **32** (2012) 1091–1100.
- [29] F. Merlevède, M. Peligrad and S. Utev. Recent advances in invariance principles for stationary sequences. *Probab. Surv.* **3** (2006) 1–36.
- [30] C. A. Morales, M. J. Pacifico and E. R. Pujals. Robust transitive singular sets for 3-flows are partially hyperbolic attractors or repellers. *Ann. of Math.* **160** (2004) 375–432.
- [31] I. I. Ovsyannikov and D. V. Turaev. Analytic proof of the existence of the Lorenz attractor in the extended Lorenz model. *Nonlinearity* **30** (2017) 115–137.
- [32] E. Rio. *Théorie asymptotique des processus aléatoires faiblement dépendants*. Mathématiques & Applications (Berlin) [Mathematics & Applications] **31**, Springer-Verlag, Berlin, 2000.
- [33] R. J. Serfling. Moment inequalities for the maximum cumulative sum. *Ann. Math. Statist.* **41** (1970) 1227–1234.
- [34] S. Smale. Differentiable dynamical systems. *Bull. Amer. Math. Soc.* **73** (1967) 747–817.

- [35] W. Tucker. A rigorous ODE solver and Smale's 14th problem. *Found. Comput. Math.* **2** (2002) 53–117.
- [36] L.-S. Young. Statistical properties of dynamical systems with some hyperbolicity. *Ann. of Math.* **147** (1998) 585–650.
- [37] L.-S. Young. Recurrence times and rates of mixing. *Israel J. Math.* **110** (1999) 153–188.