

Nonexistence of spectral gaps in Hölder spaces for continuous time dynamical systems

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Abstract

We show that there is a natural restriction on the smoothness of spaces where the transfer operator for a continuous dynamical system has a spectral gap. Such a space cannot be embedded in a Hölder space with Hölder exponent greater than $\frac{1}{2}$ unless it consists entirely of coboundaries.

1 Introduction

Decay of correlations (rates of mixing) and strong statistical properties are well-understood for Axiom A diffeomorphisms since the work of [2, 9, 10]. Mixing rates are computed with respect to any equilibrium measure with Hölder potential. Up to a finite cycle, such diffeomorphisms have exponential decay of correlations for Hölder observables. In the one-sided (uniformly expanding) setting, this is typically proved by establishing quasicompactness and a spectral gap for the associated transfer operator L . Such a spectral gap yields a decay rate $\|L^n v - \int v\| \leq C_v e^{-an}$ for v Hölder, where $\|\cdot\|$ is a suitable Hölder norm and a, C_v are positive constants. Decay of correlations for Hölder observables is an immediate consequence of the decay for L^n . This philosophy has been extended to large classes of nonuniformly expanding dynamical systems with exponential [13] and subexponential decay of correlations [14].

For continuous time dynamical systems, the usual techniques [5, 7, 8] bypass spectral gaps; the only exceptions that we know of being Tsujii [11, 12]. However, the result in [11] is for suspension semiflows over the doubling map with a C^3 roof function, where the smoothness of the roof function is crucial and very restrictive. A spectral gap for contact Anosov flows is obtained in [12]; unfortunately it seems nontrivial to extend this to nonuniformly hyperbolic contact flows (or uniformly hyperbolic

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contact flows with unbounded distortion), see [1] which proves exponential decay of correlations for billiard flows with a contact structure but does not establish a spectral gap. Indeed, apart from [11, 12], there are no results on spectral gaps of transfer operators for semiflows and flows.

The results of Tsujii [11, 12] provide a spectral gap in an anisotropic Banach space. In this paper we obtain a restriction on the Banach spaces that can yield a spectral gap. We work in the following general setting:

Let (Λ, d) be a bounded metric space with Borel probability measure μ , and let $F_t : \Lambda \rightarrow \Lambda$ be a measure-preserving semiflow. We suppose that $t \rightarrow F_t$ is Lipschitz a.e. on Λ . Let $L_t : L^1(\Lambda) \rightarrow L^1(\Lambda)$ denote the transfer operator corresponding to F_t (so $\int_{\Lambda} L_t v w d\mu = \int_{\Lambda} v w \circ F_t d\mu$ for all $v \in L^1(\Lambda)$, $w \in L^\infty(\Lambda)$, $t > 0$). Let $v \in L^\infty(\Lambda)$ and define $v_t = \int_0^t v \circ F_r dr$ for $t \geq 0$.

Theorem 1.1 *Let $\eta \in (\frac{1}{2}, 1)$. Suppose that $L_t v \in C^\eta(\Lambda)$ for all $t > 0$ and that $\int_0^\infty \|L_t v\|_\eta dt < \infty$. Then v_t is a coboundary:*

$$v_t = \chi \circ F_t - \chi \quad \text{for all } t \geq 0, \text{ a.e. on } \Lambda$$

where $\chi = \int_0^\infty L_t v dt \in C^\eta(\Lambda)$. In particular, $\sup_{t \geq 0} |v_t|_\infty < \infty$.

Here, $|g|_\infty = \text{ess sup}_\Lambda |g|$ and $\|g\|_\eta = |g|_\infty + \sup_{x \neq y} |g(x) - g(y)|/d(x, y)^\eta$.

Theorem 1.1 implies that any Banach space admitting a spectral gap and embedded in $C^\eta(\Lambda)$ for some $\eta > \frac{1}{2}$ is cohomologically trivial. However, for (non)uniformly expanding semiflows and (non)uniformly hyperbolic flows of the type in the aforementioned references, coboundaries are known to be exceedingly rare, see for example [3, Section 2.3.3]. Hence, Theorem 1.1 can be viewed as an ‘‘anti-spectral gap’’ result for such continuous time dynamical systems.

2 Proof of Theorem 1.1

Let $v \in L^\infty(\Lambda)$, with $L_t v \in C^\eta(\Lambda)$ for all $t > 0$ and $\int_0^\infty \|L_t v\|_\eta dt < \infty$ where $\eta \in (\frac{1}{2}, 1)$. Following Gordin [6] we consider a martingale-coboundary decomposition. Define $\chi = \int_0^\infty L_t v dt \in C^\eta(\Lambda)$, and

$$v_t = \int_0^t v \circ F_r dr, \quad m_t = v_t - \chi \circ F_t + \chi,$$

for $t \geq 0$. Let \mathcal{B} denote the Borel σ -algebra on Λ .

Proposition 2.1 *(i) $t \rightarrow m_t$ is C^η a.e. on Λ .*

(ii) $\mathbb{E}(m_t | F_t^{-1} \mathcal{B}) = 0$ for all $t \geq 0$.

Proof (i) For $0 \leq s \leq t \leq 1$ and $x \in \Lambda$,

$$\begin{aligned} |m_s(x) - m_t(x)| &\leq |v_s(x) - v_t(x)| + |\chi(F_s x) - \chi(F_t x)| \\ &\leq |s - t| \|v\|_\infty + |\chi|_\eta d(F_s x, F_t x)^\eta. \end{aligned}$$

Since $t \mapsto F_t$ is a.e. Lipschitz, it follows that $t \mapsto m_t$ is a.e. C^η .

(ii) Let $U_t v = v \circ F_t$, and recall that $L_t U_t = I$ and $\mathbb{E}(\cdot | F_t^{-1} \mathcal{B}) = U_t L_t$. Then

$$\begin{aligned} L_t m_t &= L_t(v_t - U_t \chi + \chi) = \int_0^t L_t U_r v dr - \chi + \int_0^\infty L_t L_r v dr \\ &= \int_0^t L_{t-r} v dr - \chi + \int_0^\infty L_{t+r} v dr = \int_0^t L_r v dr - \chi + \int_t^\infty L_r v dr = 0. \end{aligned}$$

Hence $\mathbb{E}(m_t | F_t^{-1} \mathcal{B}) = U_t L_t m_t = 0$. ■

Proof of Theorem 1.1 Fix $T > 0$, and define

$$M_T(t) = m_T - m_{T-t} = m_t \circ F_{T-t}, \quad t \in [0, T].$$

Also, define the filtration $\mathcal{G}_{T,t} = F_{T-t}^{-1} \mathcal{B}$. It is immediate that $M_T(t) = m_t \circ F_{T-t}$ is $\mathcal{G}_{T,t}$ -measurable. Also, for $s < t$ we have $M_T(t) - M_T(s) = m_{T-s} - m_{T-t} = m_{t-s} \circ F_{T-t}$, so

$$\mathbb{E}(M_T(t) - M_T(s) | \mathcal{G}_{T,s}) = \mathbb{E}(m_{t-s} \circ F_{T-t} | F_{T-s}^{-1} \mathcal{B}) = \mathbb{E}(m_{t-s} | F_{t-s}^{-1} \mathcal{B}) \circ F_{T-t} = 0$$

by Proposition 2.1(ii). Hence M_T is a martingale for each $T > 0$. Next,

$$|M_T(t)|_\infty = |m_t \circ F_{T-t}|_\infty \leq |m_t|_\infty \leq |v_t|_\infty + 2|\chi|_\infty \leq T|v|_\infty + 2|\chi|_\infty.$$

Hence $M_T(t)$, $t \in [0, T]$, is a bounded martingale.

By Proposition 2.1(i), M_T has C^η sample paths. Since $\eta > \frac{1}{2}$, it follows from general martingale theory that $M_T \equiv 0$ a.e. Taking $t = T$, we obtain $m_T = 0$ a.e. Hence $v_T = \chi \circ F_T - \chi$ a.e. for all $T > 0$ as required.

For completeness, we include the argument that $M_T \equiv 0$ a.e. We require two standard properties of the quadratic variation process $t \mapsto [M_T](t)$; a reference for these is [4, Theorem 4.1]. First, $[M_T](t)$ is the limit in probability as $n \rightarrow \infty$ of

$$S_n(t) = \sum_{j=1}^n (M_T(jt/n) - M_T((j-1)t/n))^2.$$

Second (noting that $M_T(0) = 0$),

$$[M_T](t) = M_T(t)^2 - 2 \int_0^t M_T dM_T,$$

where the stochastic integral has expectation zero. In particular, $\mathbb{E}([M_T]) \equiv \mathbb{E}(M_T^2)$.

Since M_T has Hölder sample paths with exponent $\eta > \frac{1}{2}$, we have a.e. that

$$|S_n(t)| = O(t^\eta n^{-(2\eta-1)}) \rightarrow 0 \quad \text{as } n \rightarrow \infty.$$

Hence $[M_T] \equiv 0$ a.e. It follows that $\mathbb{E}(M_T^2) \equiv 0$ and so $M_T \equiv 0$ a.e. ■

References

- [1] V. Baladi, M. F. Demers and C. Liverani. Exponential decay of correlations for finite horizon Sinai billiard flows. *Invent. Math.* **211** (2018) 39–177.
- [2] R. Bowen. *Equilibrium States and the Ergodic Theory of Anosov Diffeomorphisms*. Lecture Notes in Math. **470**, Springer, Berlin, 1975.
- [3] I. Chevyrev, P. K. Friz, A. Korepanov, I. Melbourne and H. Zhang. Multiscale systems, homogenization, and rough paths. *Probability and Analysis in Interacting Physical Systems: In Honor of S.R.S. Varadhan, Berlin, August, 2016*” (P. Friz et al., ed.), Springer Proceedings in Mathematics & Statistics **283**, Springer, 2019, pp. 17–48.
- [4] K. L. Chung and R. J. Williams. *Introduction to stochastic integration*, second ed., Probability and its Applications, Birkhäuser Boston, Inc., Boston, MA, 1990.
- [5] D. Dolgopyat. On the decay of correlations in Anosov flows. *Ann. of Math.* **147** (1998) 357–390.
- [6] M. I. Gordin. The central limit theorem for stationary processes. *Soviet Math. Dokl.* **10** (1969) 1174–1176.
- [7] C. Liverani. On contact Anosov flows. *Ann. of Math.* **159** (2004) 1275–1312.
- [8] M. Pollicott. On the rate of mixing of Axiom A flows. *Invent. Math.* **81** (1985) 413–426.
- [9] D. Ruelle. *Thermodynamic Formalism*. Encyclopedia of Math. and its Applications **5**, Addison Wesley, Massachusetts, 1978.
- [10] Y. G. Sinai. Gibbs measures in ergodic theory. *Russ. Math. Surv.* **27** (1972) 21–70.
- [11] M. Tsujii. Decay of correlations in suspension semi-flows of angle-multiplying maps. *Ergodic Theory Dynam. Systems* **28** (2008) 291–317.
- [12] M. Tsujii. Quasi-compactness of transfer operators for contact Anosov flows. *Nonlinearity* **23** (2010) 1495–1545.
- [13] L.-S. Young. Statistical properties of dynamical systems with some hyperbolicity. *Ann. of Math.* **147** (1998) 585–650.
- [14] L.-S. Young. Recurrence times and rates of mixing. *Israel J. Math.* **110** (1999) 153–188.